

Dynamic Aperture/Dynamic Momentum Aperture



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- *How the transverse single particle dynamics limits the dynamic aperture (injection efficiency) and the momentum aperture and therefore the beam lifetime in Light Sources (also applicable to colliders, damping rings, ...)*

Topics



- ❑ Motivation/Introduction
- ❑ Dynamic Aperture/Injection Efficiency
 - ❑ Machine Model Calibration
 - ❑ Frequency Map Analysis
 - ❑ Measurements of Frequency Maps
- ❑ Dynamic Momentum Aperture
 - ❑ RF-Amplitude Scan
 - ❑ Off-Energy Frequency Maps
- ❑ Summary

I will discuss three different examples (one injection efficiency and two Touschek lifetime examples), where the insight gained by frequency map analysis was applied directly to significantly improve the performance of the ALS.

Motivation



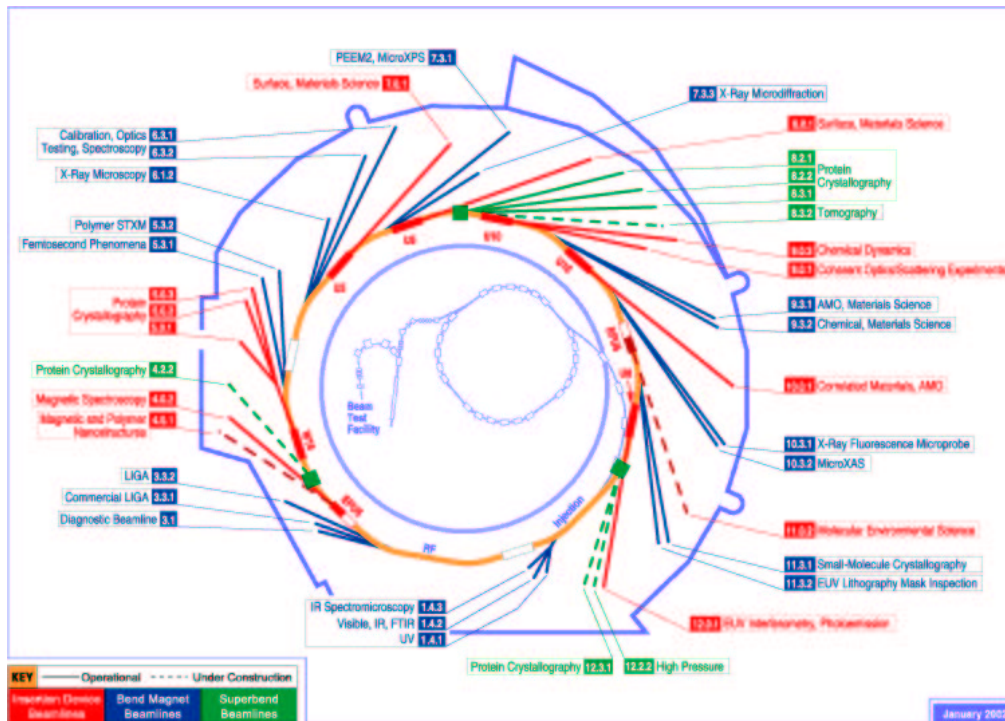
One performance limitation of ALS (like for other 3rd generation synchrotron radiation sources) is **single particle dynamics**

- ❑ Beam lifetime (**dynamic momentum aperture**)
- ❑ Injection efficiency (**dynamic aperture**)
- ❑ Transverse beam profile (**resonant tail generation**)

Gaining a good model (ideal machine with random errors is insufficient) and understanding of beam **dynamics at high amplitudes** can provide

- ❑ **Diagnostic** of problems
- ❑ **Improvement** of actual performance
- ❑ **Prediction** what will happen due to future expansions (superbends, distributed dispersion, lower beta functions, mini gap IDs, femtosecond)

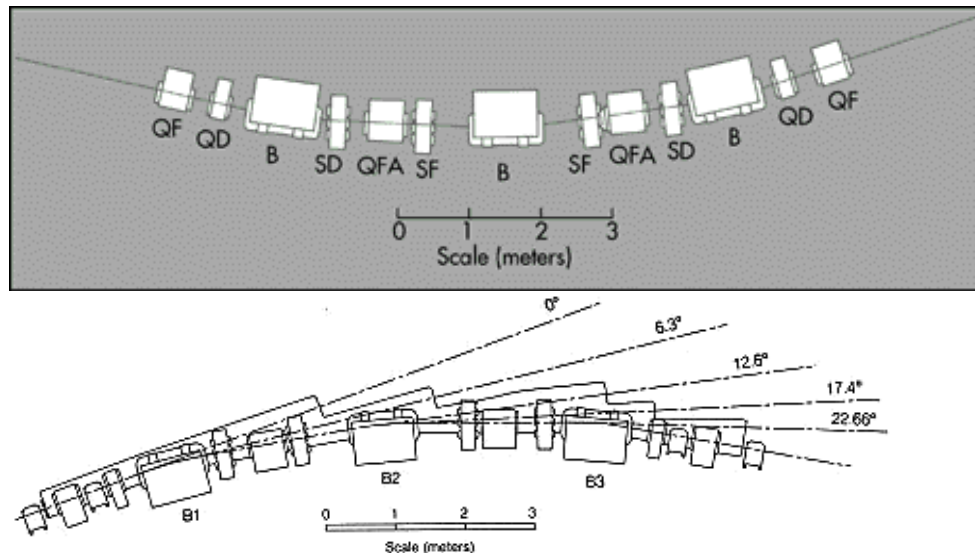
ALS Parameters



Nominal energy	1.5–1.9 GeV
Circumference	196.8 m
RF frequency	500 MHz
Harmonic number	328
Revolution freq.	1.5 MHz
Bunch current	1–2 mA or 5–35 mA
Energy spread	$6-8 (8-10) \times 10^{-4}$
Bunch length	15 – 20 ps

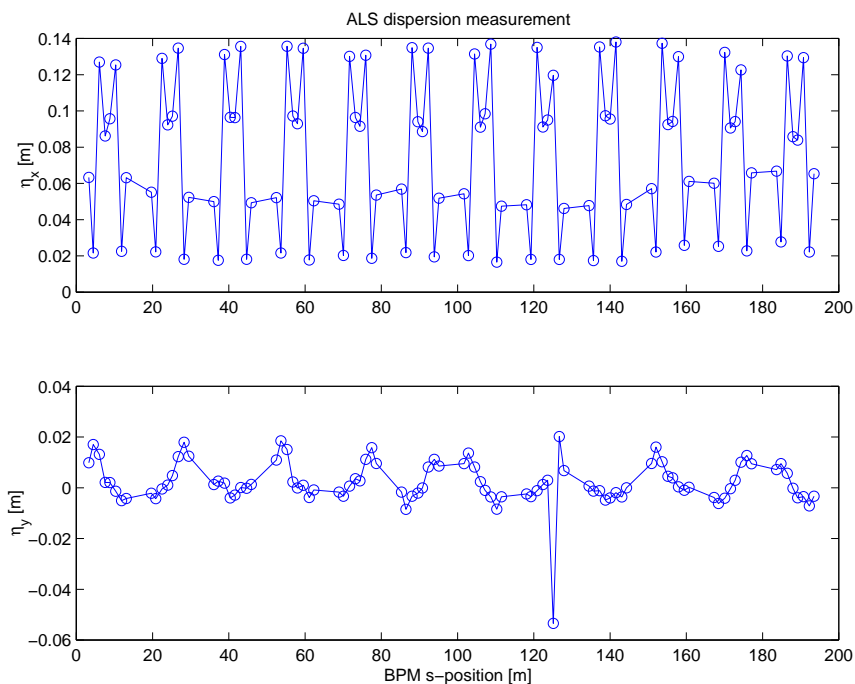
$1 \text{ mA} = \frac{2}{3} \text{ nC} = 4.2 \times 10^9 \text{ electrons}$
 About 1 day a week available for accelerator physics studies (most studies directly related to user operation)

The ALS Lattice



- 12 nearly identical arcs → TBA
- 8 fast beam position monitors per arc, 1024 turn recording capability; additional slow straight section IDBPMs
- 8 horizontal, 6 vertical corrector magnets per arc
- no individual skew quadrupoles so far, only 4 families
- TBA lattice very insensitive to dipole kicks (vertical dispersion)
- x-ray diagnostic beamline on bending magnet

Vertical Dispersion/Coupling



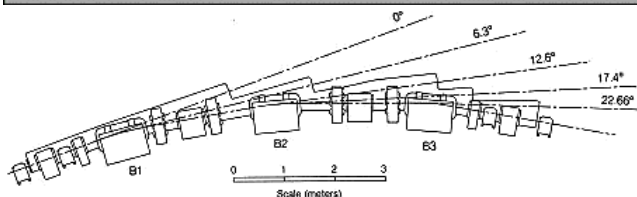
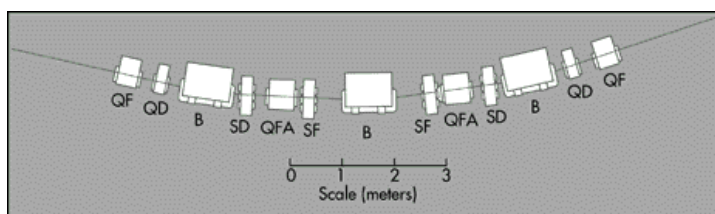
24-Jun-2002

- With normal orbit correction (BBA) spurious vertical dispersion typically 5 mm rms, with coupling correction about 2 mm rms
- Betatron coupling small \Rightarrow vertical emittance dominated by dispersion (at 1.5 GeV $\frac{\varepsilon_y}{\varepsilon_x}$ routinely below 1%, achieved values below 0.1% (5 pm rad))
- With Superbends (2001) change to distributed dispersion lattice

Response Matrix/Machine Model



- fitting a machine model to the response matrix (SVD, LOCO)

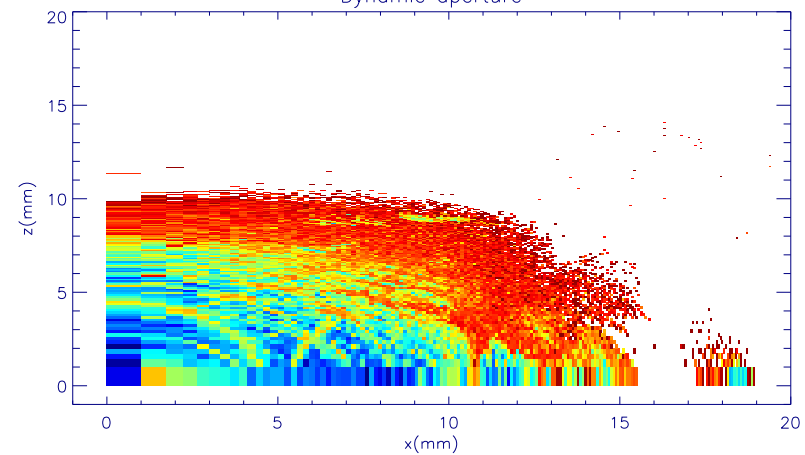
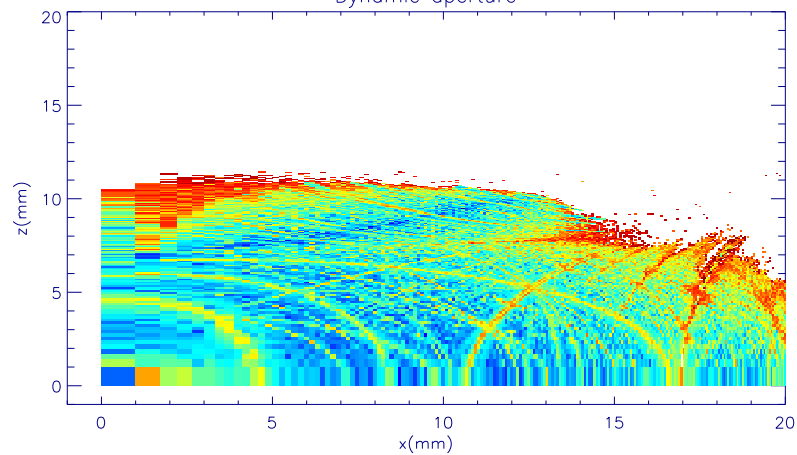
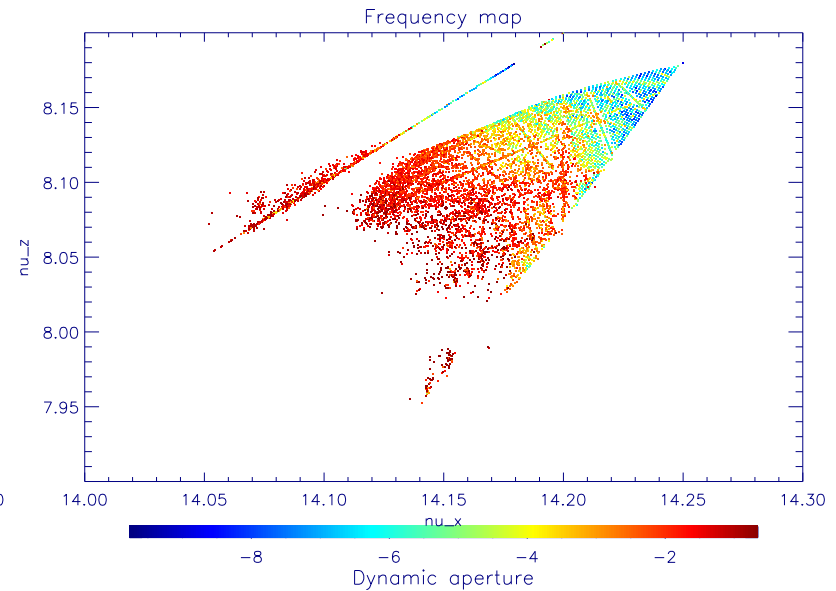
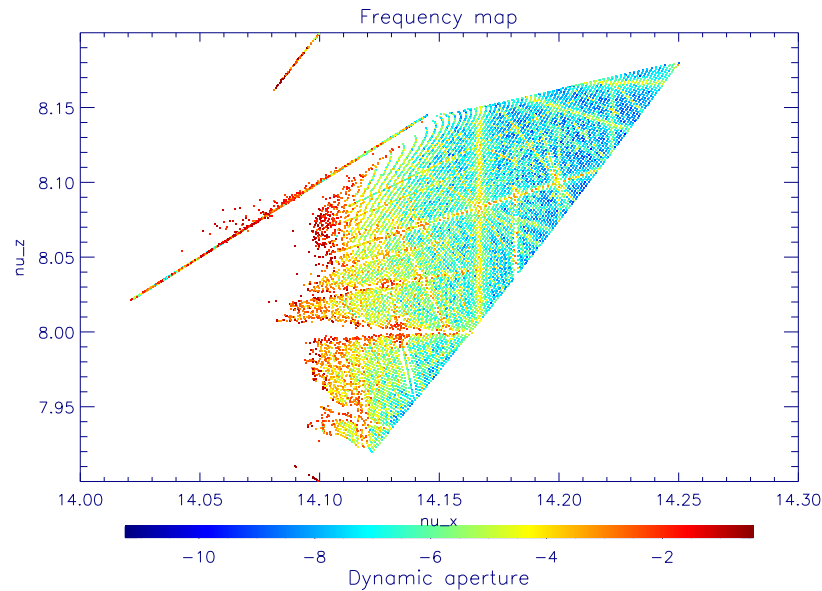


$$C_{12}^{ij} = \left[R^{ij} (1 - R^{jj})^{-1} \right]_{12} - \frac{\eta_i \eta_j}{\left(\alpha - \frac{1}{\gamma^2} \right) C}$$

$$\hat{C}^{ij} = C^{ij} + \sum_k \frac{\partial C^{ij}}{\partial g_k} \delta g_k + C^{ij} \Delta x^i - C^{ij} \Delta y^j$$

- can determine individual quadrupole strengths
- can determine localized coupling strengths

ALS: Ideal Lattice/Calibrated Model



Ideal versus Calibrated Model



Linear lattice errors fundamentally change the **beam dynamics**

Ideal model:

- ◆ **Dynamic aperture** is large
- ◆ **Chaotic zones** at high amplitudes are small
- ◆ **Particle loss** is fast
- ◆ Particle loss due to allowed **high order resonances**

Calibrated model (linear errors):

- ◆ **Dynamic aperture** is smaller
- ◆ Large **chaotic zones**
- ◆ **Particle loss** is slow (diffusion)
- ◆ Particle loss due to unallowed **lower order resonances**

Is *either of these models* an accurate description of dynamics at high amplitudes in *real ring*? \Rightarrow test possible with *Measured Frequency Maps*

Experimental Procedure



Experimental Hardware

- ❑ horizontal + vertical single turn kicker
- ❑ 96 turn by turn monitors (1024 turns)

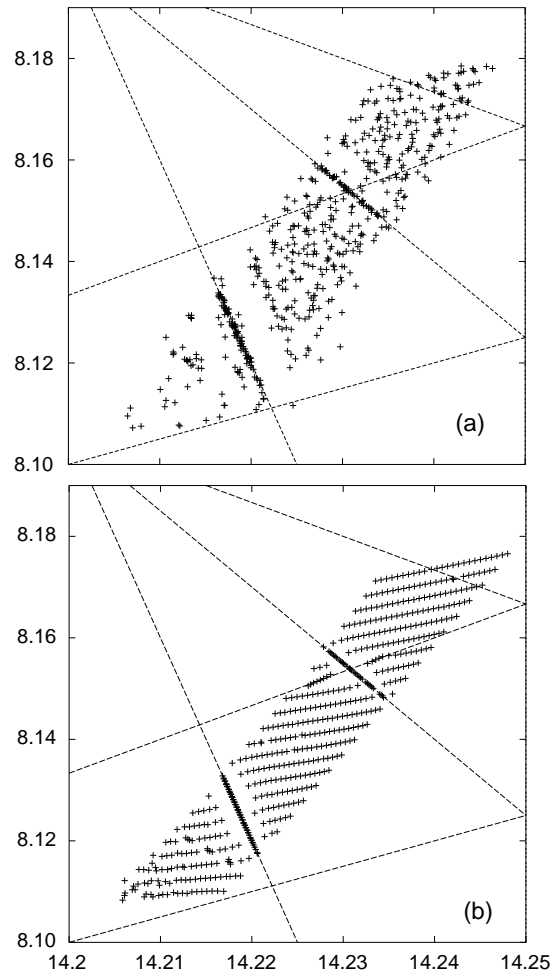
Experimental Procedure

- ❑ Electron beam (single bunch or small bunch train) gets simultaneously a horizontal and a vertical kick
- ❑ Beam centroid oscillations are recorded turn by turn for 1024 turns
- ❑ Repeat with different initial conditions (hor. + vert. kick amplitude) → 400-600 total points per map

Data Analysis

- ❑ turn by turn data is analyzed with frequency analysis post processor (NAFF) and results plotted in tune plane

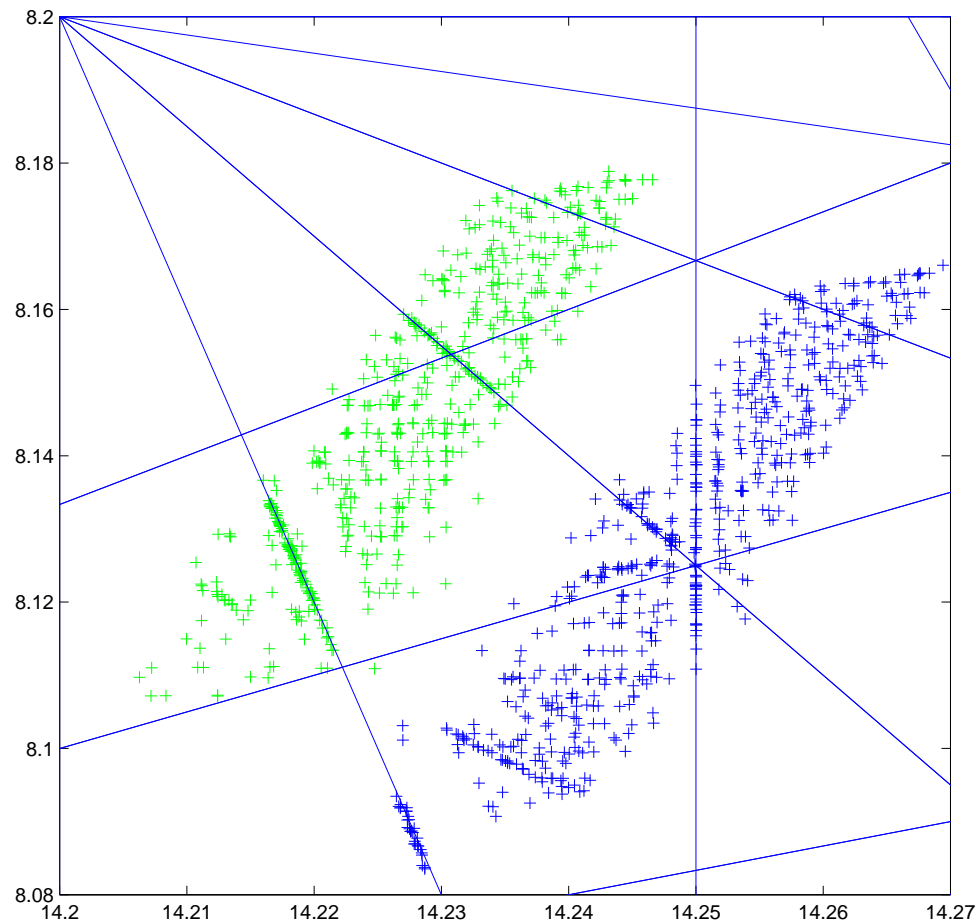
ALS: Measured Frequency Map



□ excellent agreement, using calibrated model (gradient errors), random skew errors, nominal sextupoles

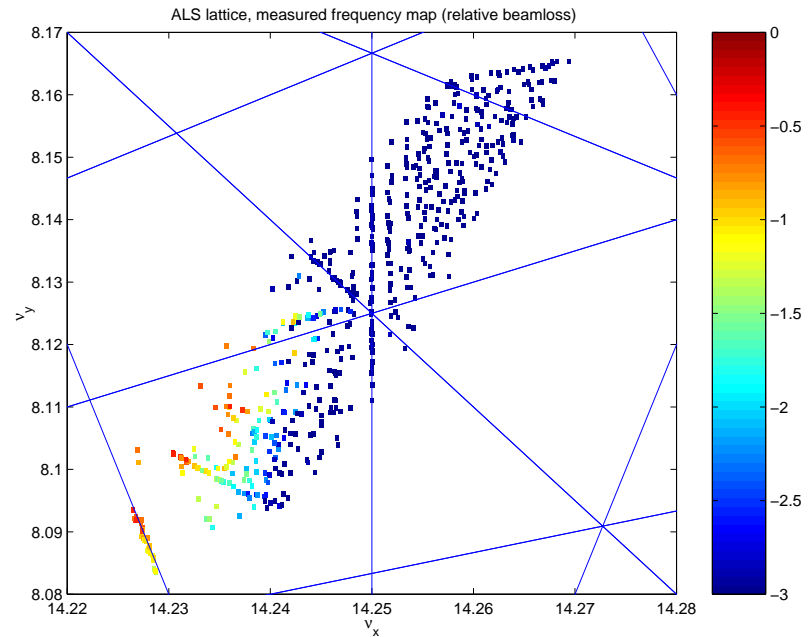
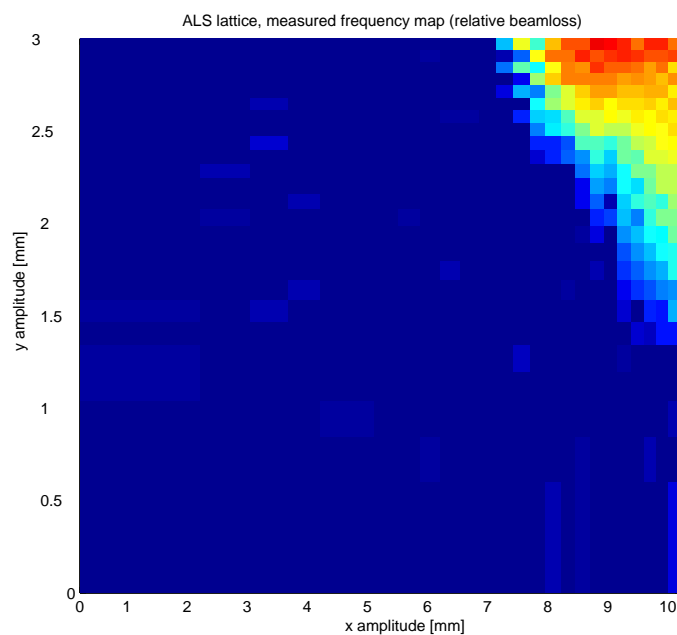
Phys. Rev. Lett. 85, 3, (July 2000), pp.558-561

ALS: Measured Frequency Maps



□ measurement → model independent diagnostic: (isolated, weak resonances + regular regions versus intersecting, strong resonances + large chaotic regions)

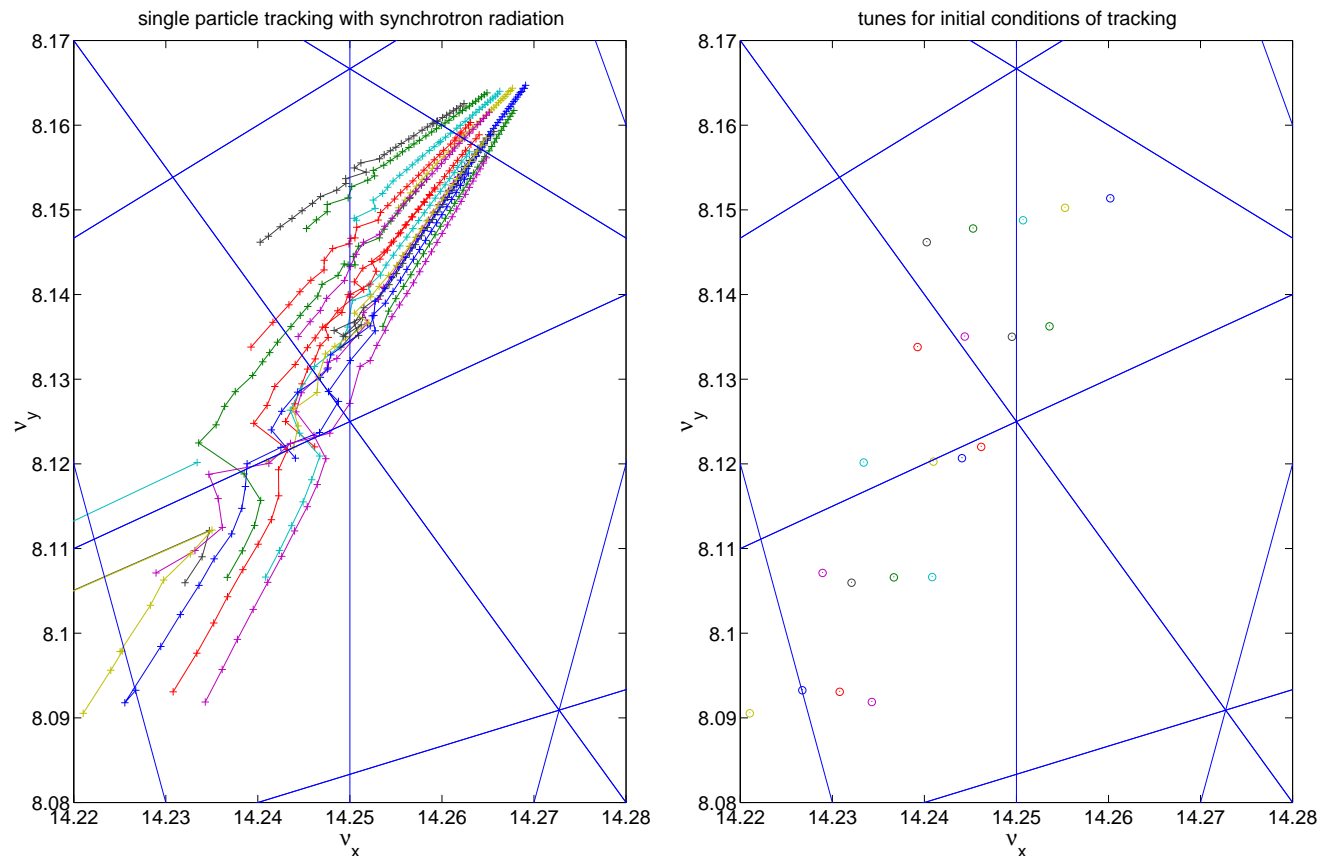
ALS: Measured Frequency Map with Measured Beam Loss



- Partial **beam loss** mostly if particles have to pass (radiation damping) through **resonance intersection**

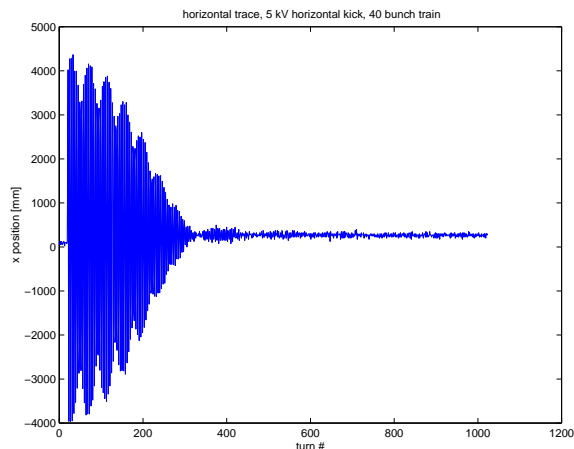
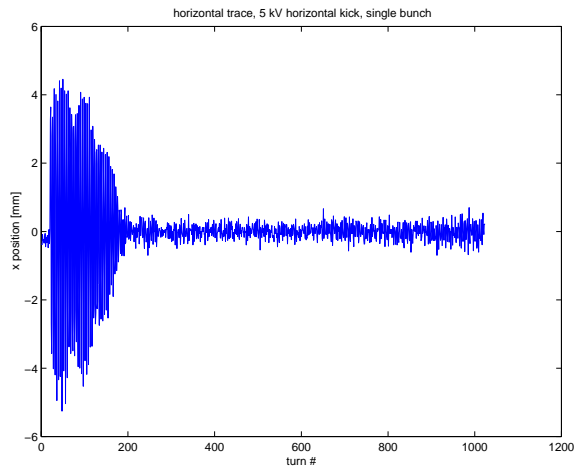
Published in *Phys. Rev. Lett.* 85, 3, (July 2000), pp.558-561

Tracking Individual Particles in Tune Space (with radiation)



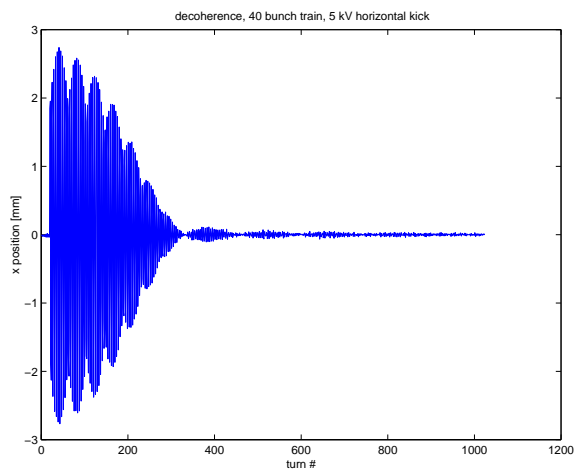
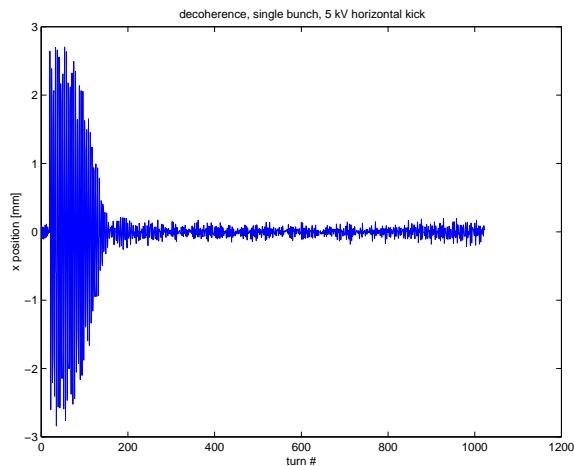
- Fraction of particles damping through resonance intersection gain enough vertical oscillation amplitude to get lost on physical vertical aperture

Experimental Difficulties: BPM noise



- **FADs** only have 8 Bit resolution
- input **Bandpass filter** has a transient step response time of about 60-80 ns \Rightarrow single bunch data very noisy
- use **40 bucket train** instead \Rightarrow possible problems: **additional decoherence, multibunch instabilities, resonance broadening**
- noise level in 40 bucket case mainly determined by digitization resolution

Decoherence

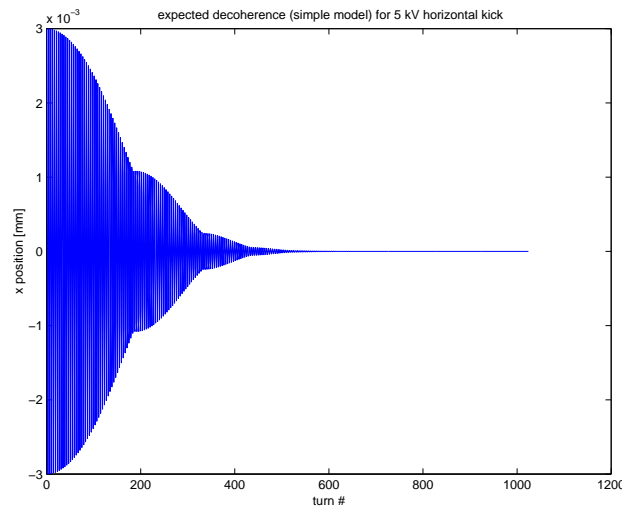


- horizontally the additional decoherence due to the bunch train is **small**
- **vertically it is significant** (kick to about same amplitude, beamsizes different by factor of 10)
- decoherence changes not only amplitude but also oscillation phase \Rightarrow signals **are not quasiperiodic** after about 10-50 turns
- one needs fast converging frequency calculation algorithms (NAFF) and low noise position data \rightarrow **2 dedicated, fast BPMs**

Decoherence II



- simple model of decoherence: detuning proportional to betatron amplitude ($\Delta\nu \propto x^2$), no cross terms \Rightarrow analytic formula



$$\begin{aligned}
 A(N) &= \frac{1}{1 + \theta^2} \exp\left(-\frac{Z^2}{2} \frac{\theta^2}{1 + \theta^2}\right) \\
 \Delta\phi(N) &= -\frac{Z^2}{2} \frac{\theta^2}{1 + \theta^2} - 2 \arctan \theta \\
 \theta &= 4\pi\mu N \\
 A_s(N) &= \exp\left(-\frac{\alpha^2}{2}\right) \\
 \alpha &= 2\sigma_s \xi \nu_s^{-1} \sin(\pi\nu_s N)
 \end{aligned}$$

- **problems:** cross terms, resonances, BPM nonlinearities

Summary (Frequency Maps)



- *For the first time **frequency maps** have been measured on a storage ring*
 - ◆ The network of **coupling resonances** is clearly visible (**Arnold Web**)
 - ◆ The observed **dynamics** is complex
 - ◆ **Model independent** information about quality of lattice
- *Remarkable **agreement** between the measured and simulated maps*
 - ◆ ⇒ Confidence in our **machine model** (orbit response matrix analysis)
- *Outlook:*
 - ◆ Reduction in **measurement time** (to about 15 minutes) for frequency map planned ⇒ diagnostic tool for **routine operation**
 - ◆ Touschek lifetime ⇒ **off energy** frequency maps and measurements
→ studies started

Motivation (Dynamic Momentum Aperture)



Laurent Nadolski, David Robin, Christoph Steier, Ying Wu (LBNL);

Winfried Decking (DESY); Jacques Laskar(IMC-CNRS)

- **Lifetime** is a **crucial performance parameter** for all light sources \Rightarrow for 3rd generation light sources lifetime dominated by **Touschek effect** \Rightarrow Touschek lifetime strong function of **momentum aperture** ε

$$\tau_{\text{tou}} \propto E^3 \frac{V_{\text{bunch}} \sigma'_x}{I_{\text{bunch}}} \varepsilon^2 f(\varepsilon, \sigma'_x, E)$$

\Rightarrow Momentum aperture ε influenced/limited by **single particle dynamics**

- Design momentum aperture for current light sources: $\geq 3\%$;
achieved: about $\geq 2\%$ \Rightarrow Need to understand current **limitations**
- For ALS: understand+minimize **impact of future expansions** (**Superbends**, distributed dispersion, lower beta functions, mini gap IDs, femtosecond)

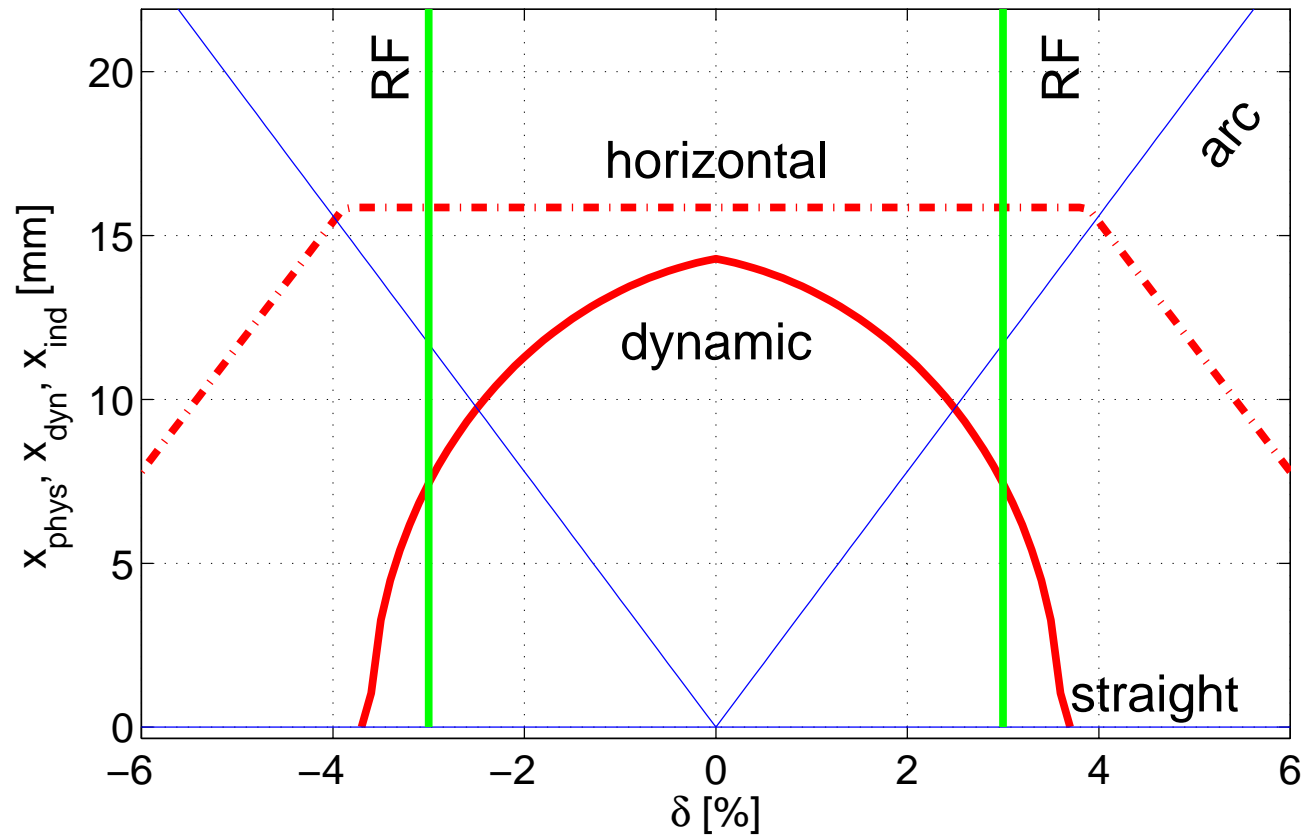
Motivation/Outline (Momentum Aperture)



- Design momentum aperture for future light sources (e.g. Soleil) 5–6% (necessary for reasonable lifetimes).
⇒ Application of some methods of **Frequency Map Analysis** can help understand the off-energy single particle dynamics.

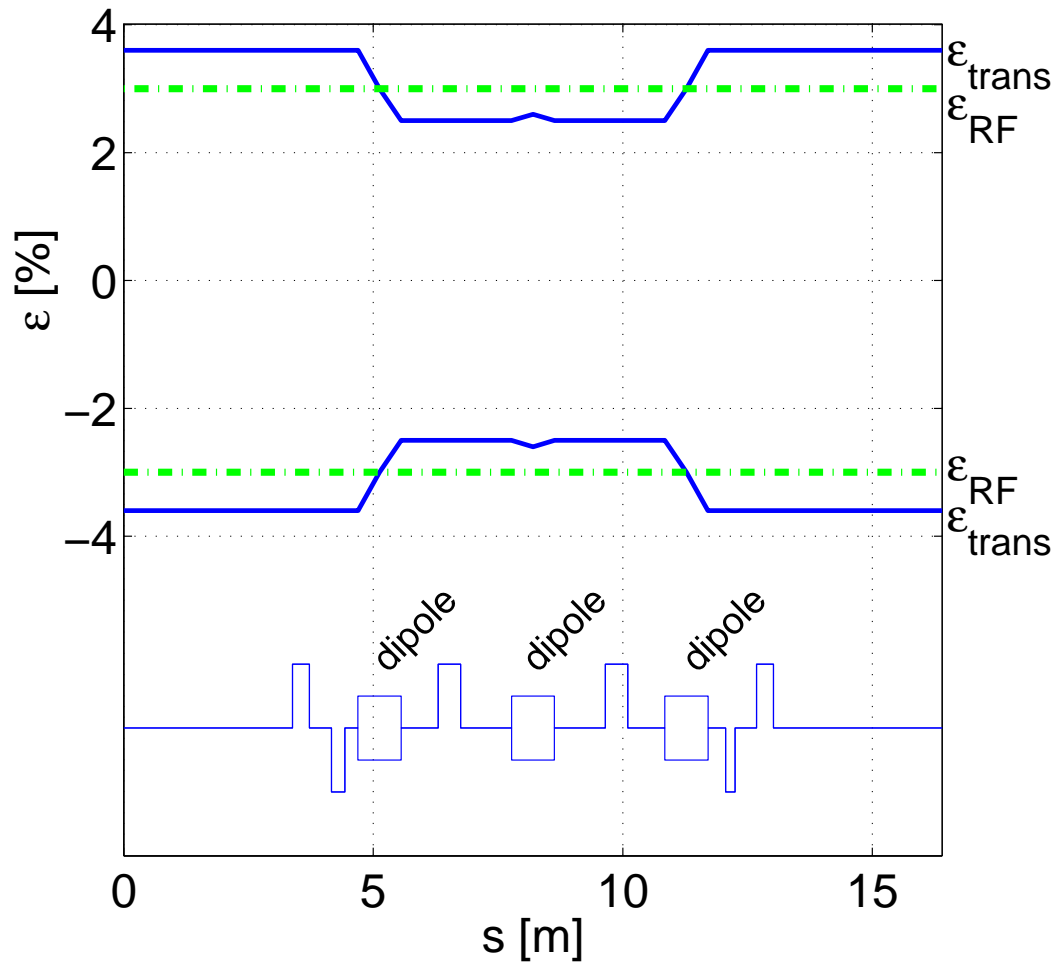
- Outline:
 - ◆ What **limits** momentum aperture
 - ◆ ALS limited by **dynamic** momentum aperture
 - ◆ Loss mechanism/**Frequency Map Analysis** (on-energy ⇒ off-energy)
 - ◆ Measurements can serve as **model independent debugging tool**
 - ◆ Agreement with simulations ⇒ tool for **predictions** for expansions/**future machines**

What Determines Momentum Aperture ?



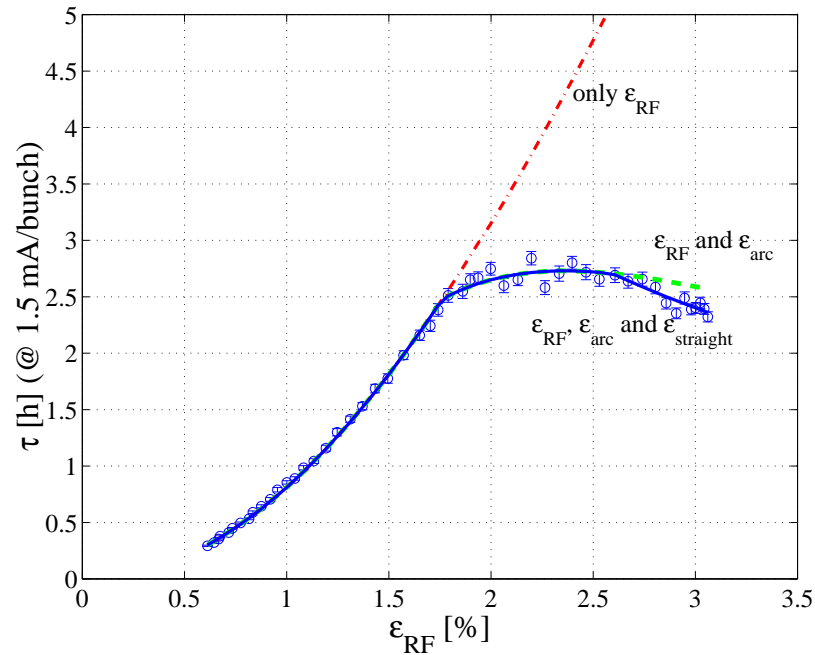
⇒ Possibility to study **momentum aperture** by **scanning rf-voltage**

Dependence of Momentum Aperture on Position in Lattice



\Rightarrow for the ALS the momentum aperture is (nearly) constant in the straights and in the arcs, respectively

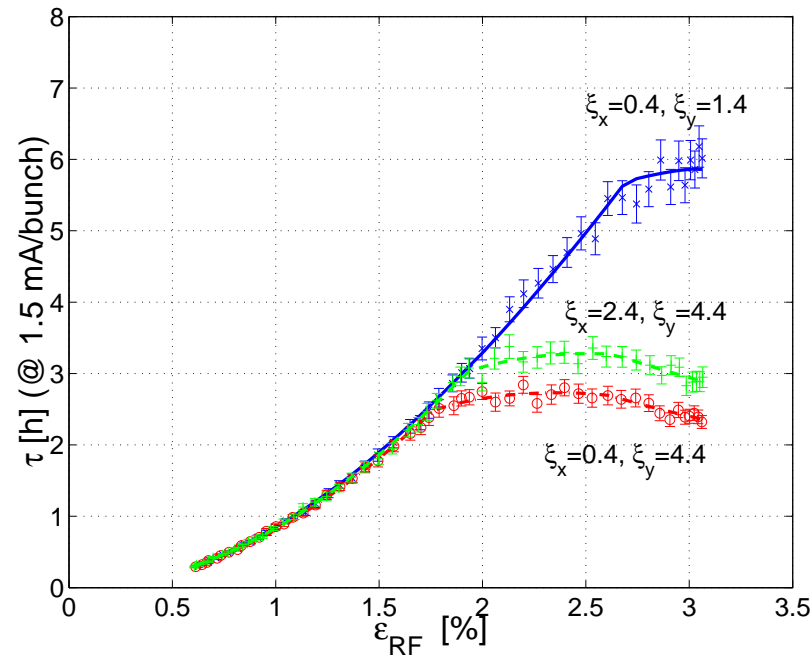
ALS: RF-amplitude scan



- Data can be fit with just three parameters:

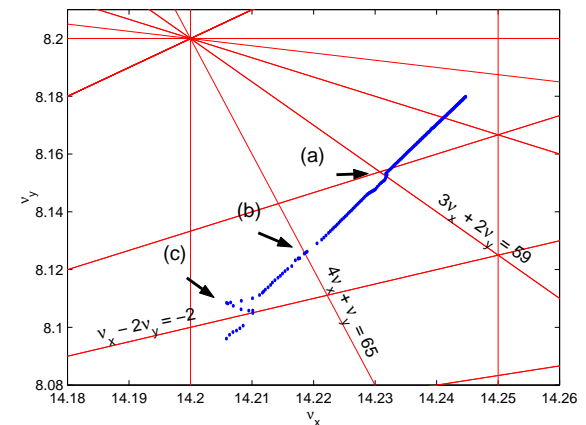
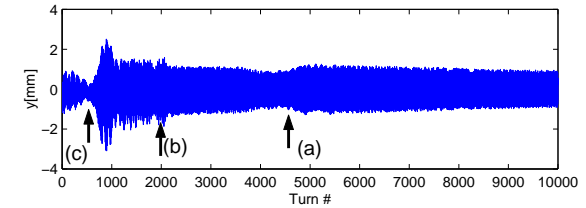
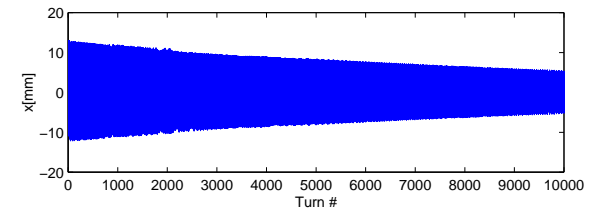
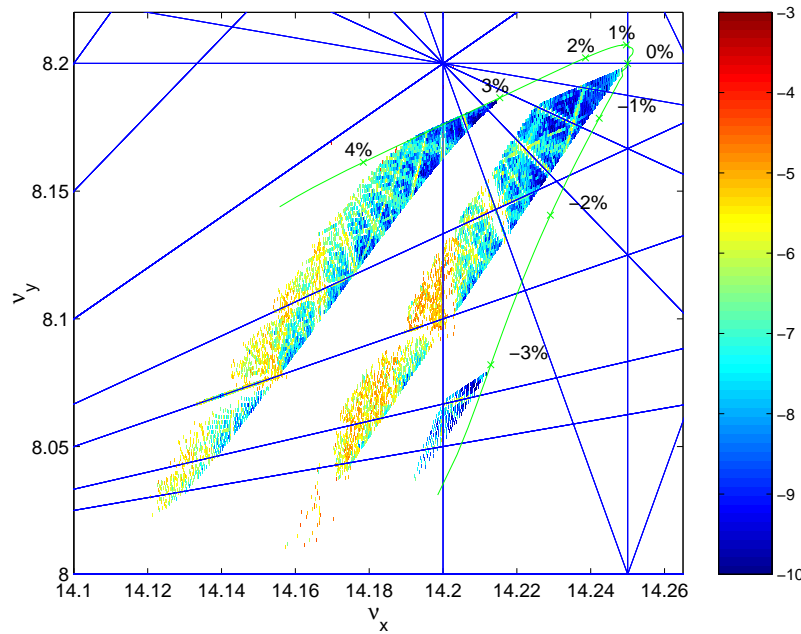
$$\tau_{tous} = \frac{AE_0^3 \sigma_L(\nu_s)}{LI_b} \left[\epsilon_{arc}^2 \int_{arc} \frac{\sigma_x(s)\sigma_y(s)\sigma'_x(s)}{\ln\left(\frac{1}{1.78} \frac{\sigma'_x(s)^2 E_0^2}{\epsilon_{arc}^2}\right)} ds + \epsilon_{straight}^2 \int_{straight} \frac{\sigma_x(s)\sigma_y(s)\sigma'_x(s)}{\ln\left(\frac{1}{1.78} \frac{\sigma'_x(s)^2 E_0^2}{\epsilon_{straight}^2}\right)} ds \right]$$

ALS: RF-amplitude scan



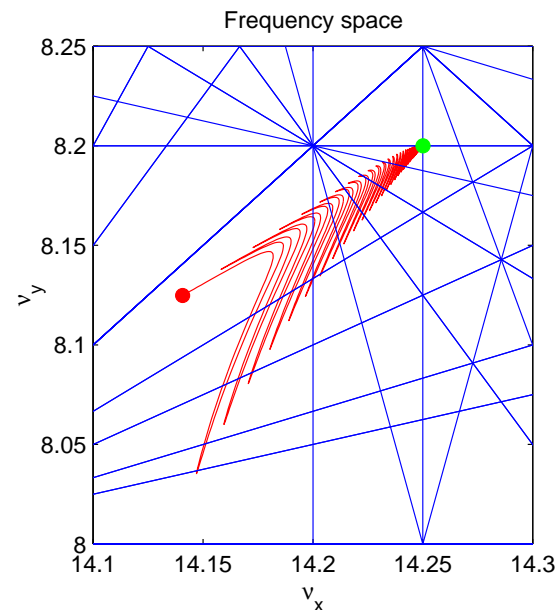
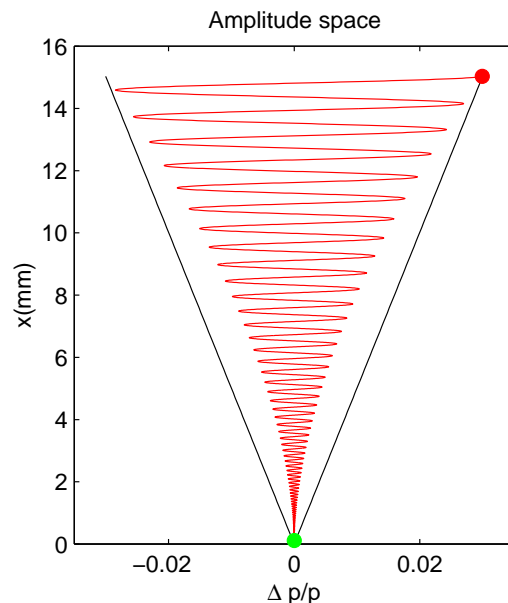
- ❑ Momentum aperture in ALS is clearly impacted by dynamics
- ❑ Sensitivity to chromaticity is (at first) surprisingly large (sextupole strength only different by few percent)

Off Momentum Frequency Map Analysis



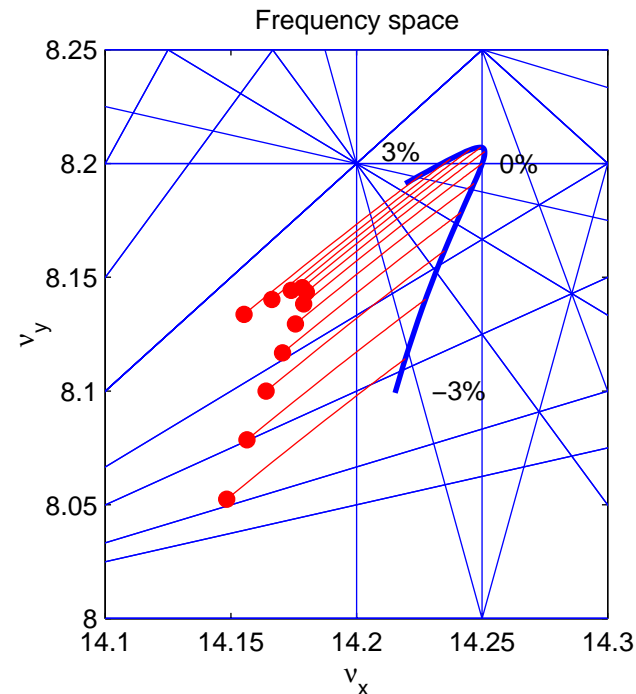
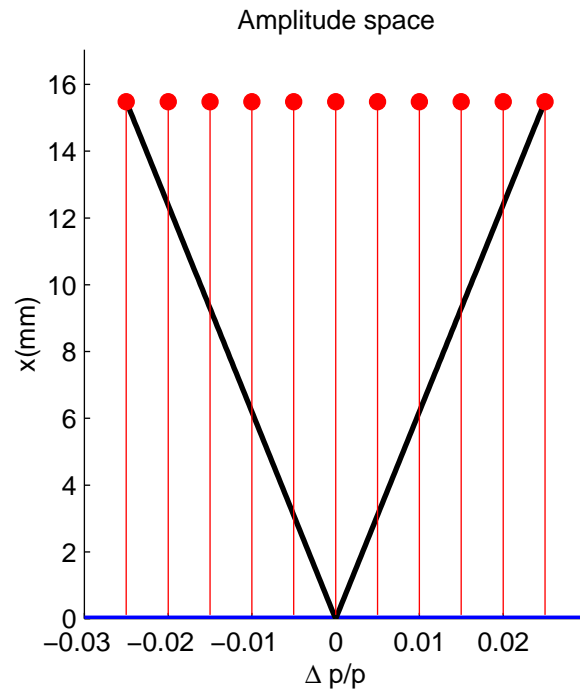
□ Knowledge of off-momentum dynamics is sufficient to understand the full six dimensional dynamics

Touschek Scattering → Tuneshift → Particle Loss



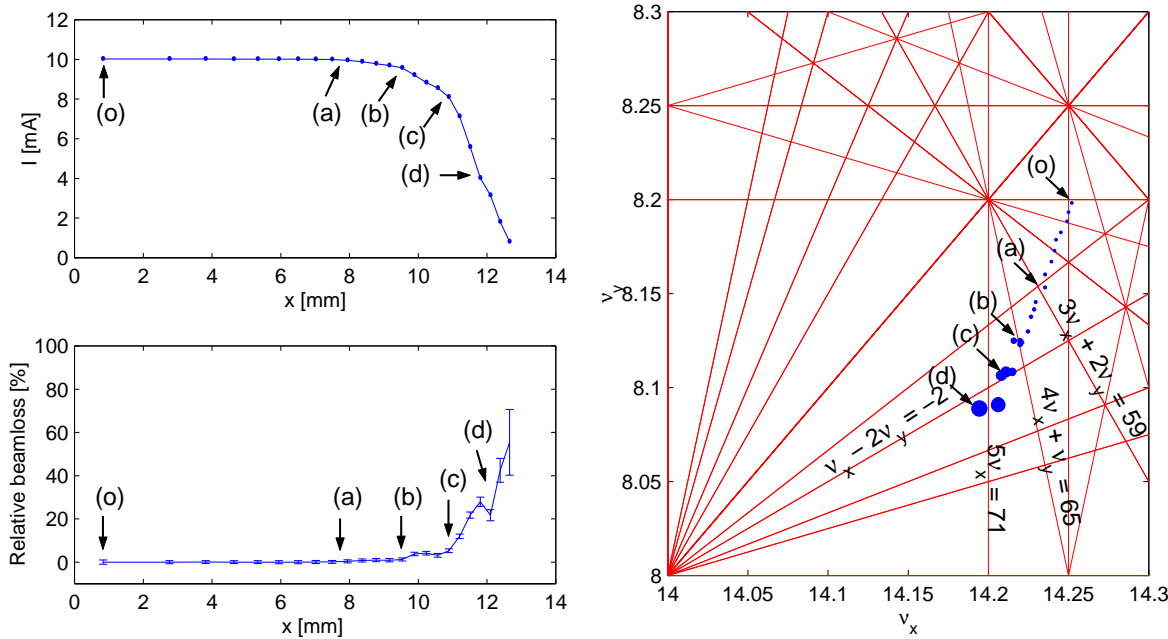
- Particle **loosing/gaining energy** → **hor.** (dispersion/ H -function) + long. oscillation
- Change in particle's **betatron tune** due to
 - ◆ Synchrotron oscillation (change in Δp , chromaticity)
 - ◆ Radiation damping ($\Delta p/p$ and A_x , chromaticity, detuning with amplitude)
 - ◆ Particle can get into regions of tune space where the **motion gets chaotic or resonantly excited**

Measurement Principle



- ❑ Experimentally **very difficult** to apply **simultaneous** transverse and longitudinal kick
- ❑ Still possible to locate loss regions when **scanning only transverse amplitude** while keeping **energy offset constant**

Measurement Method

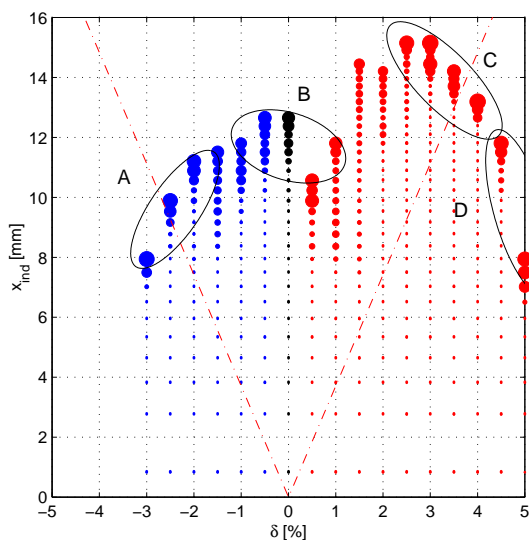


- Use pinger magnet to kick/excite the beam with increasing amplitude
- Use DCCT to record relative beam loss after each kick
- Use turn-by-turn BPMs to record oscillation frequencies

Aperture Scan – Three Different Chromaticities

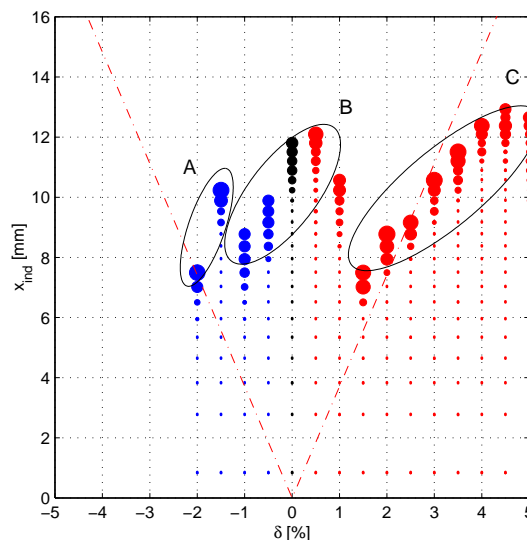


**Small Horizontal +
Small Vertical
Chromaticity**



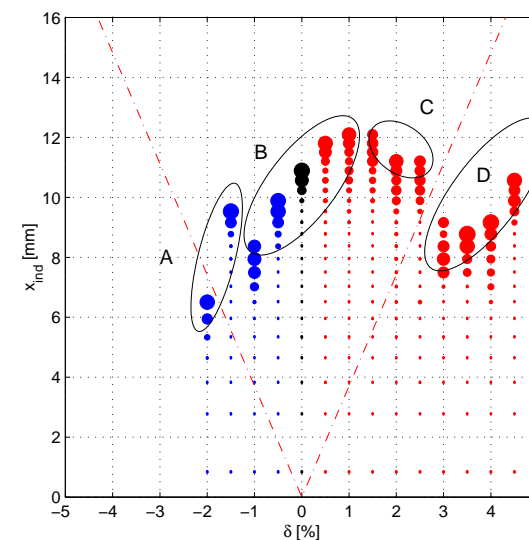
$\epsilon > 3\%$ in the straights
2.65% in the arcs

**Small Horizontal +
Large Vertical
Chromaticity**



$\epsilon = 2.6\%$ in the straights
1.75% in the arcs

**Large Horizontal +
Large Vertical
Chromaticity**



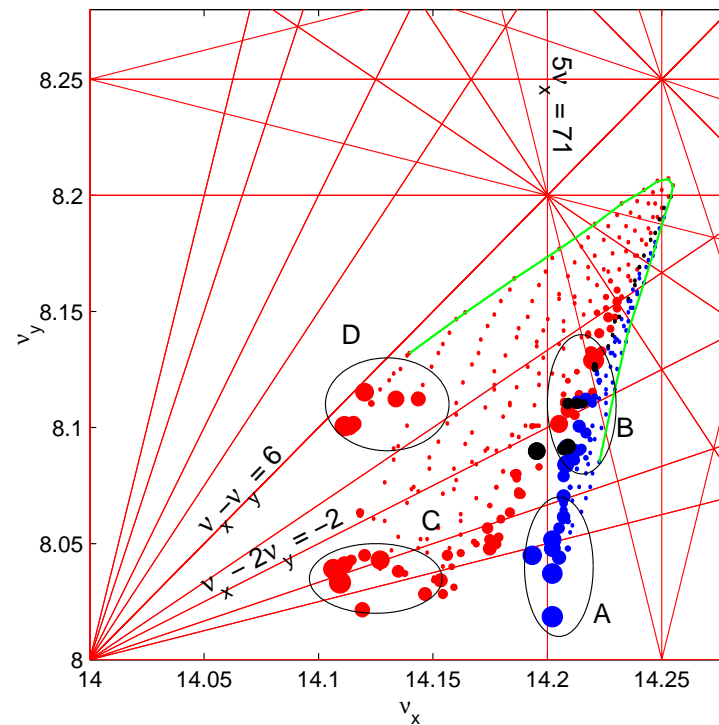
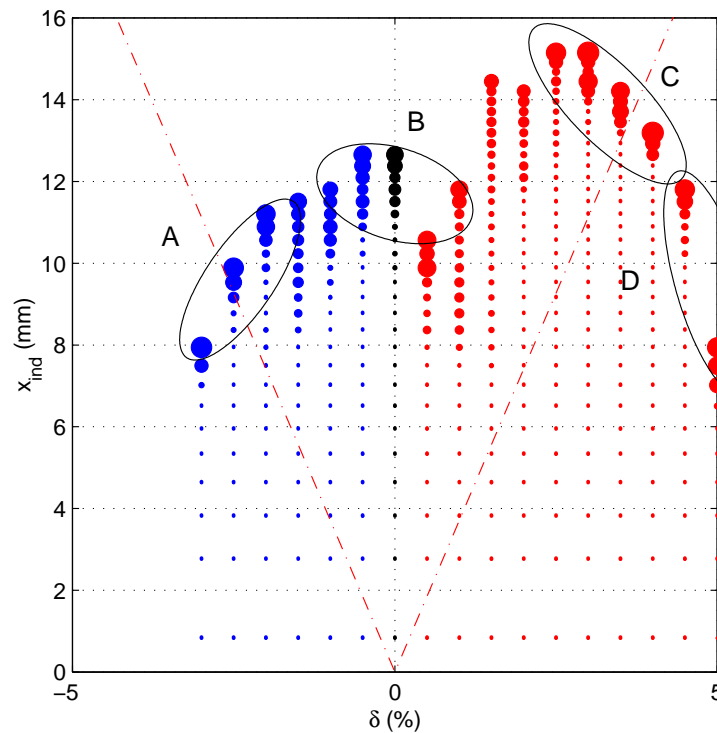
$\epsilon = 2.6\%$ in the straights
1.9% in the arcs

Experimental Difficulties: Damping Partition



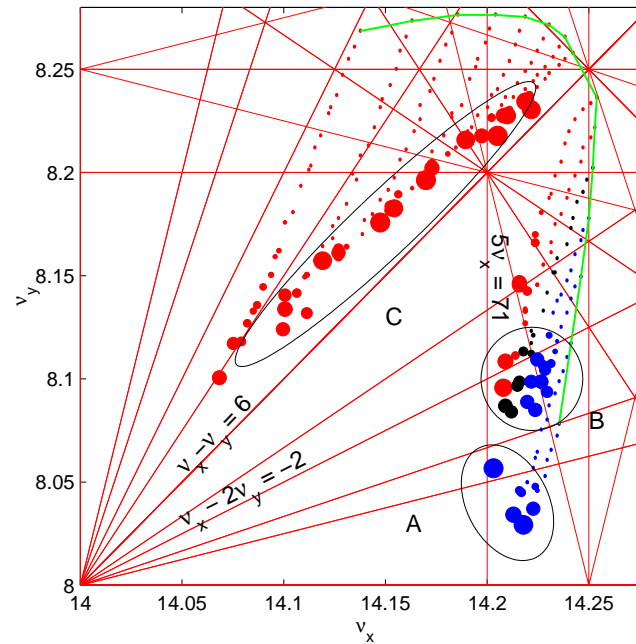
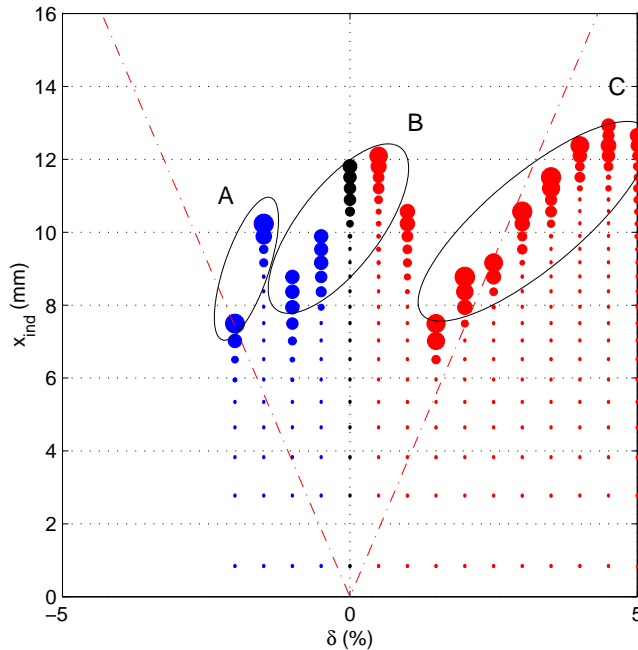
- ❑ Changing the RF-frequency to statically change the center beam energy is quite different from a particle undergoing synchrotron oscillations.
- ❑ The combination of quadrupole plus dipole field in all quadrupoles if the frequency is changed from the central frequency changes the damping partition numbers.
- ❑ The magnitude of change (for a given percentage energy change) depends on the size of the machine.
- ❑ For large machines a moderate beam energy change (changing the rf-frequency) can make the beam antidamped in one plane (typically horizontal).
- ❑ This limits how useful this (directly applied) method is for large accelerators (APS for example).

Aperture Scan – Small Chromaticity Case



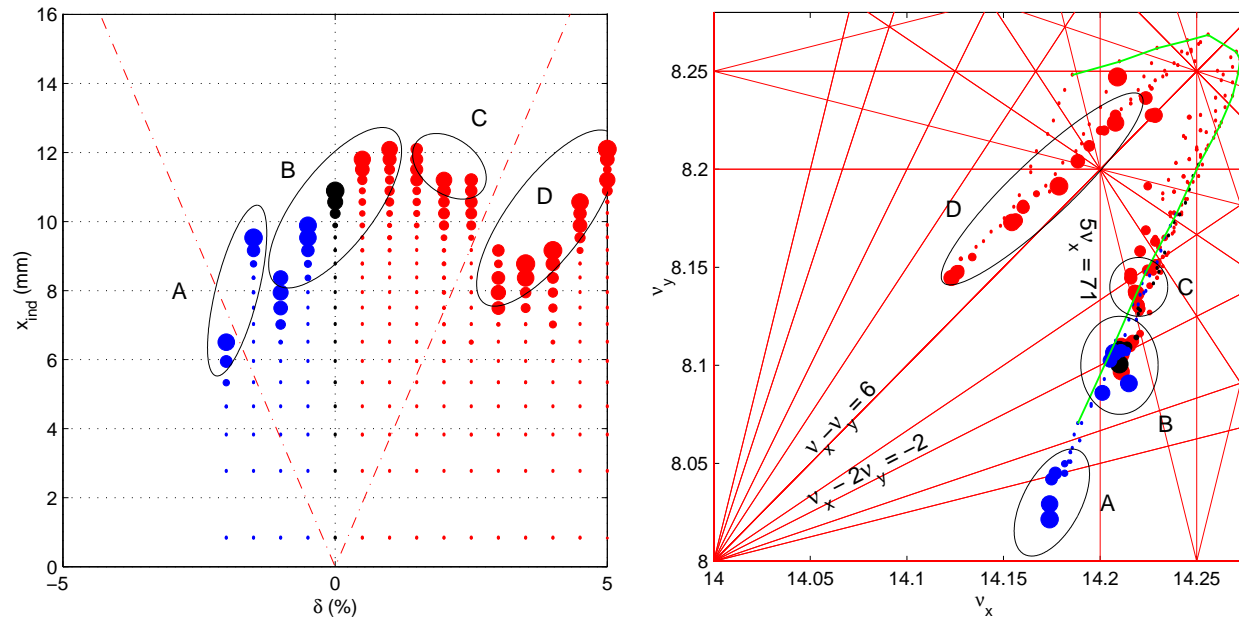
- Measured momentum aperture: $> 3\%$ (straight), 2.65% (arc)
- Ultimate limitation in both energy directions is vertical integer resonance

Large Vertical Chromaticity Case



- ❑ Measured **momentum aperture**: 2.6% (straight), 1.75% (arc)
- ❑ Limitation in negative direction is **vertical integer resonance**
- ❑ Limitation in positive direction is **coupling resonance** (used to control vertical beamsize at ALS)

Large Horizontal+Vertical Chromaticity Case

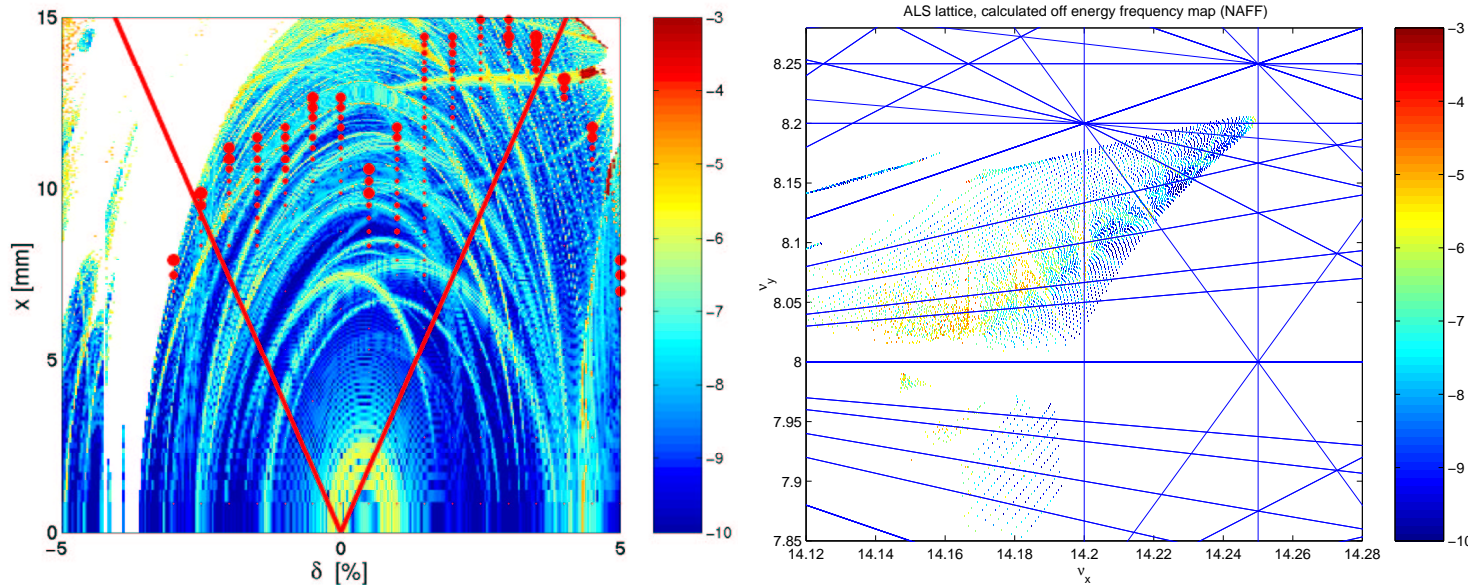


- ❑ Measured momentum aperture: 2.6% (straight), 1.9% (arc)
- ❑ Limitation in negative direction is still vertical integer resonance
- ❑ Limitation in positive direction is coupling resonance \Rightarrow shifted to larger energy deviation by additionally raising horizontal chromaticity

Comparison: Experiment - Simulation



Small Horizontal + Small Vertical Chromaticity

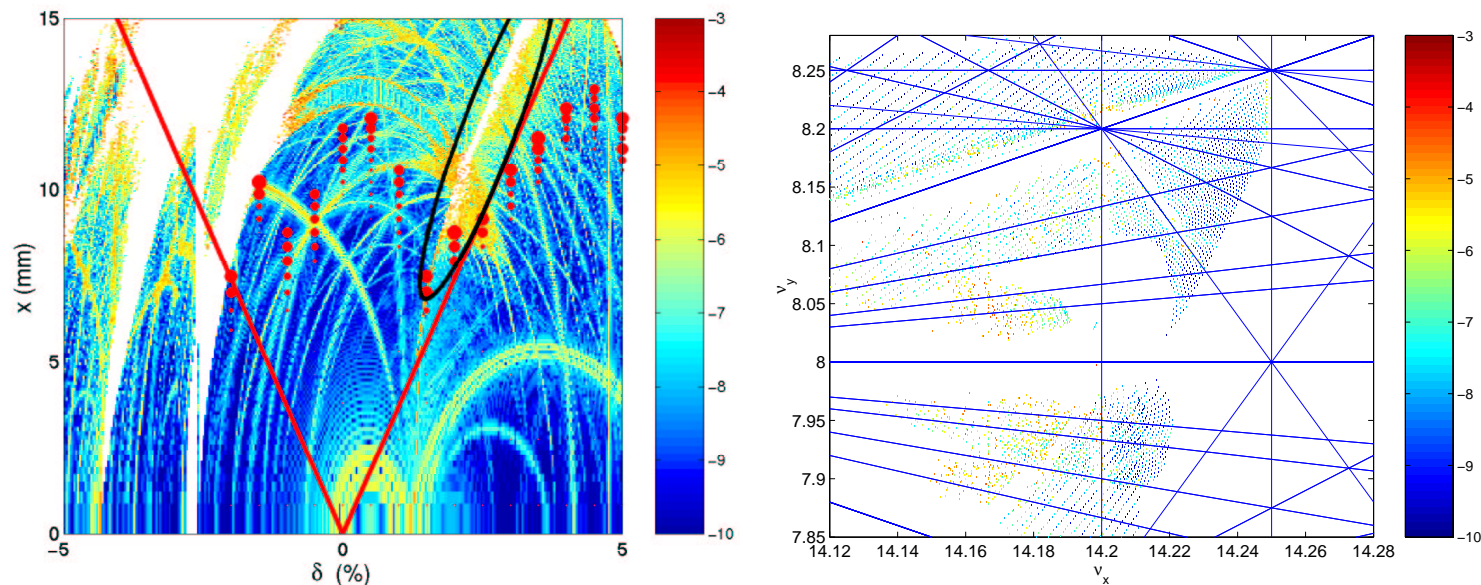


Agreement with simulations is good (though interpretation is not trivial and additional tools are necessary - plots in frequency space; comparison with experiments; off energy frequency maps, ...) \Rightarrow predict impact of new projects at ALS and performance of new light sources (and help to optimize them)

Comparison: Experiment - Simulation



Small Horizontal + Large Vertical Chromaticity

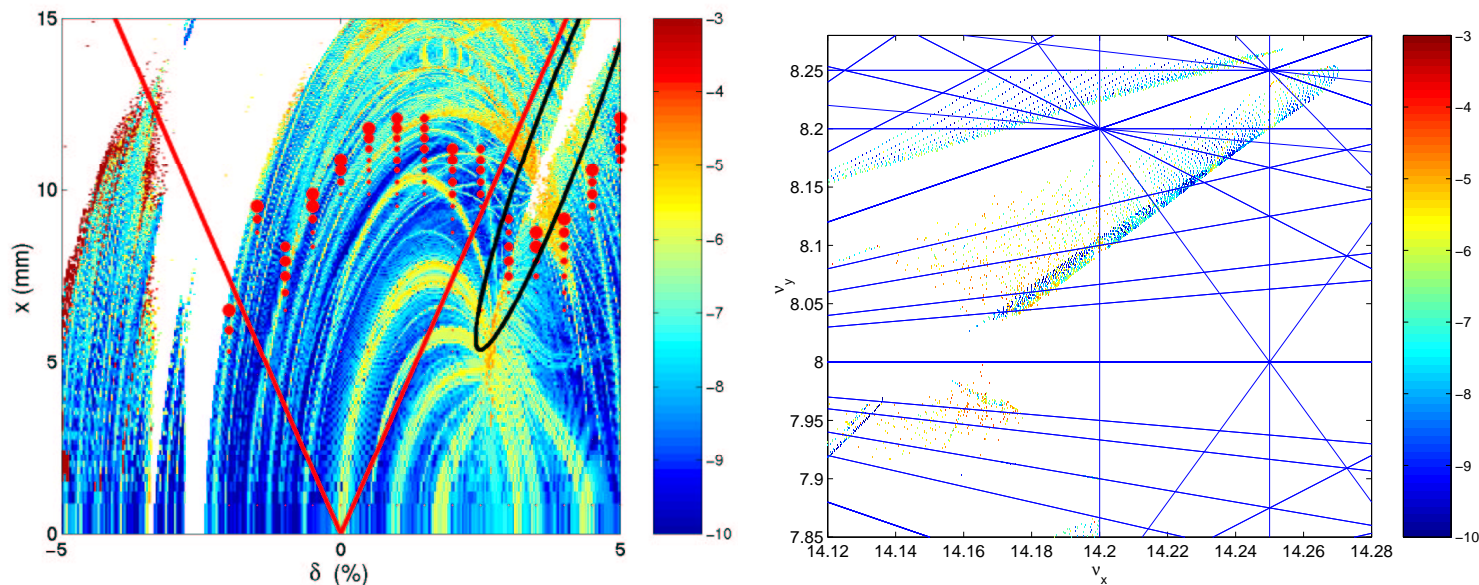


Simulations clearly reproduce reduction in momentum aperture due to coupling resonance (at positive delta) and integer resonance (at negative)

Comparison: Experiment - Simulation



Large Horizontal + Large Vertical Chromaticity

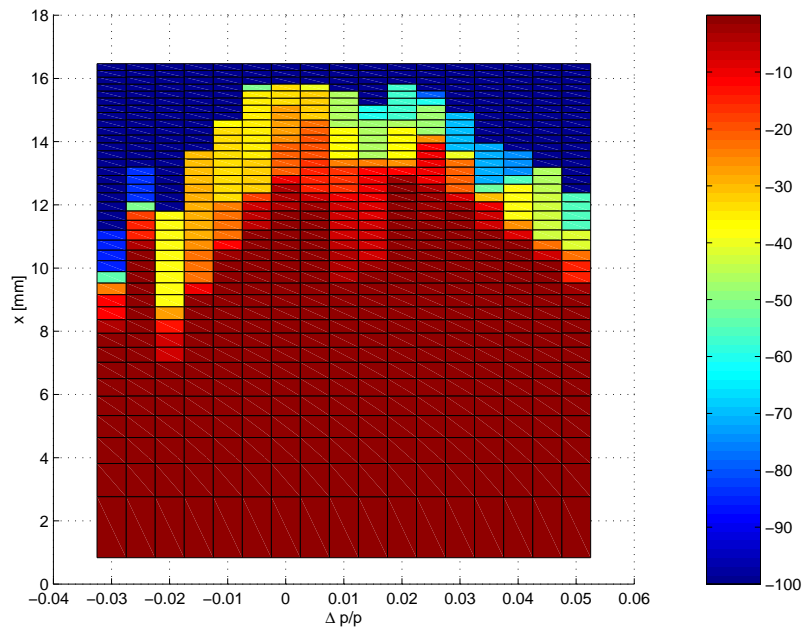
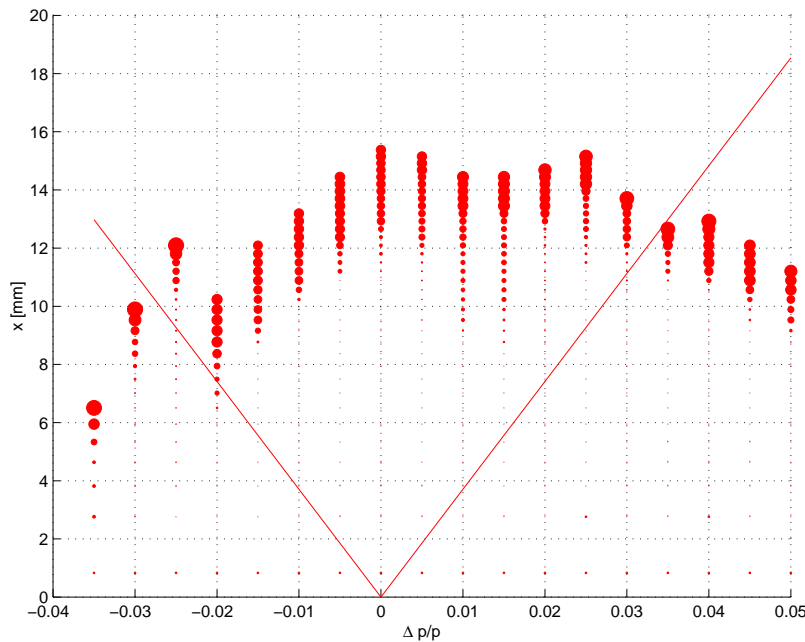


Simulations show shift of beam loss area due to coupling resonance to higher positive momentum deviations.

Off energy aperture studies/frequency maps

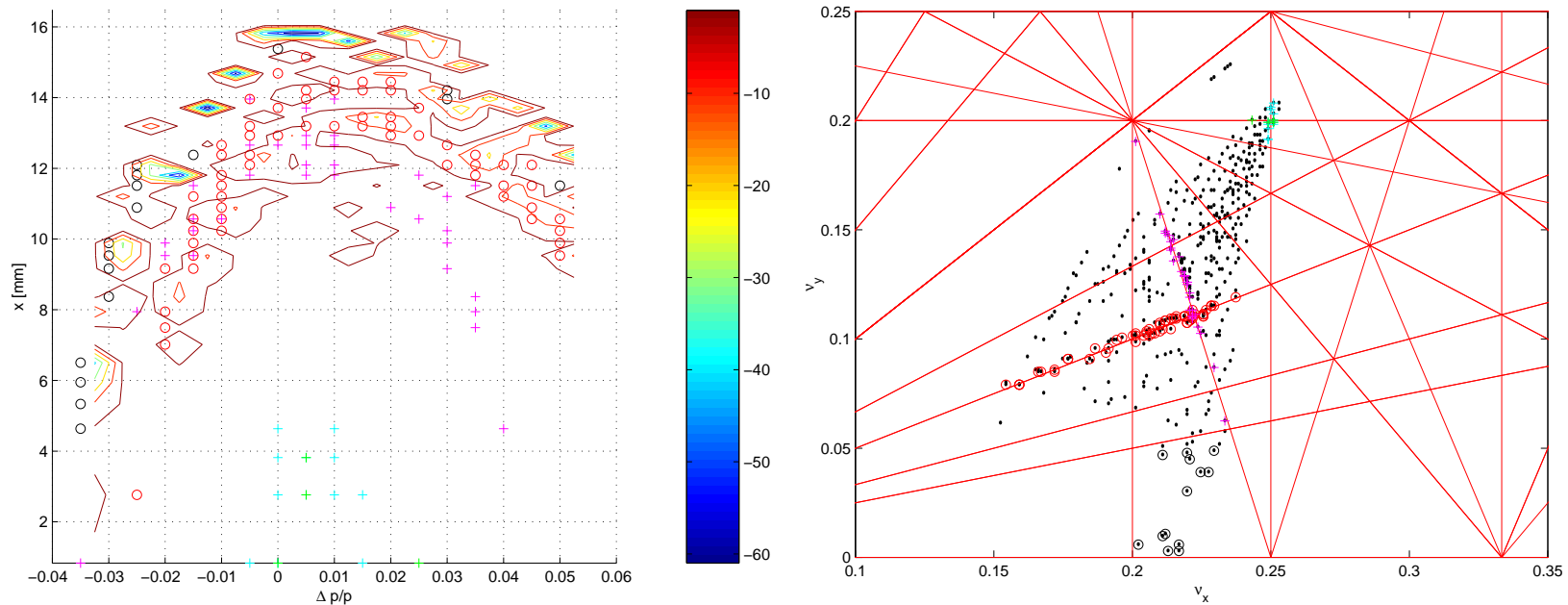


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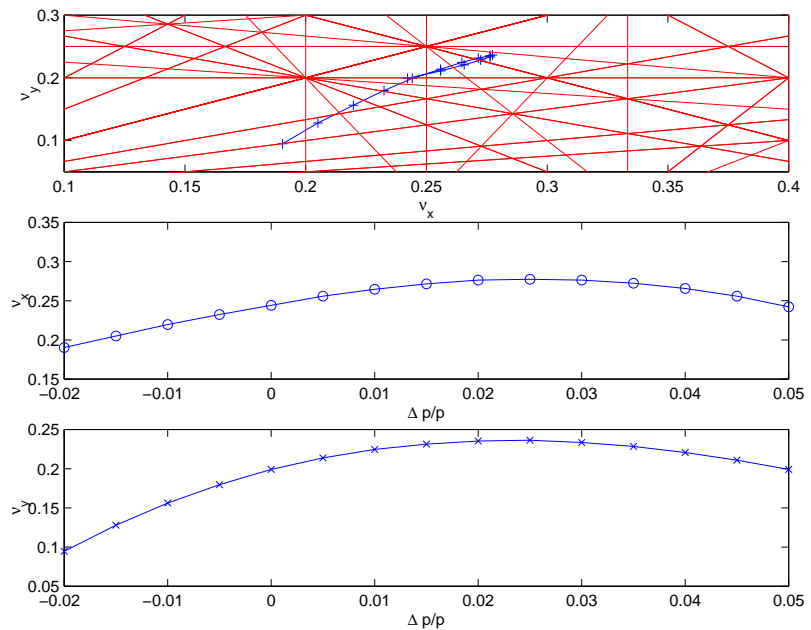
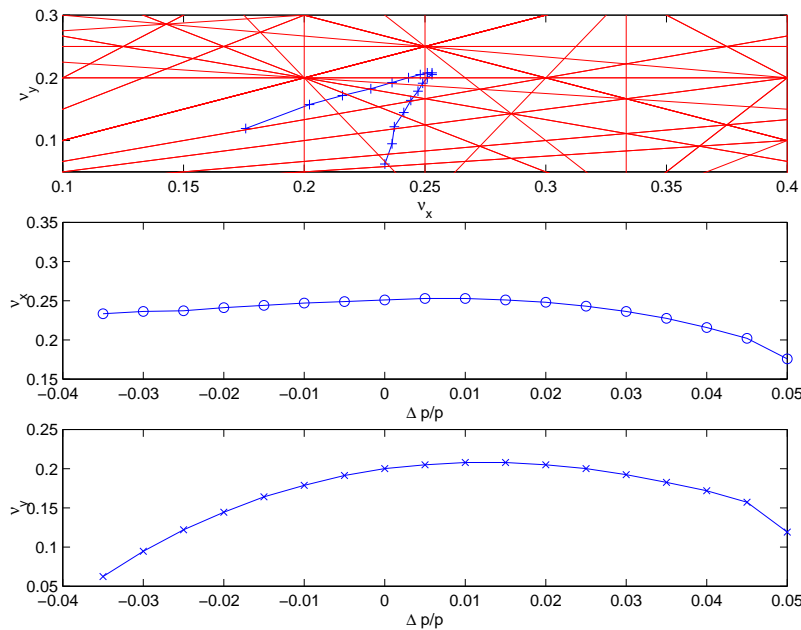
- again some regions with partial **beam loss** \Rightarrow **resonances**
- short term dynamic aperture is **not a smooth** function of energy deviation

Off energy aperture studies/frequency maps II



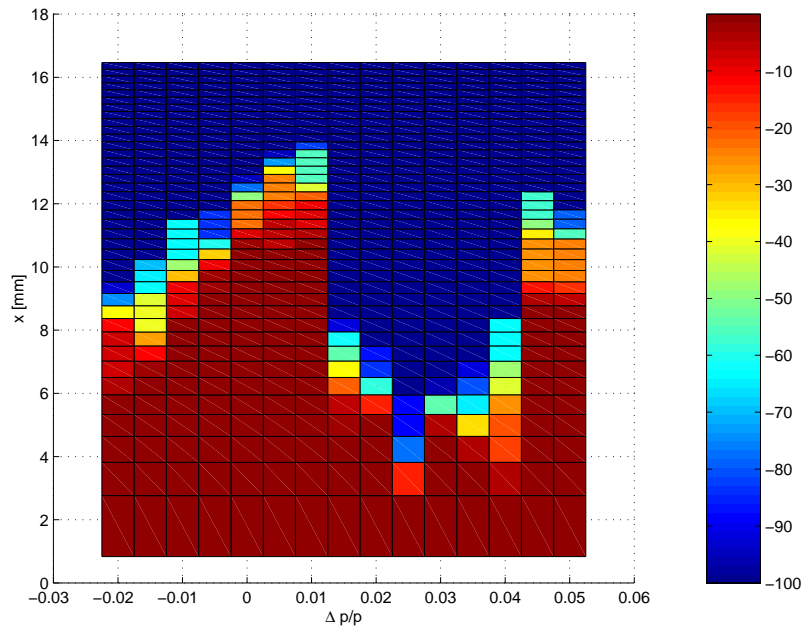
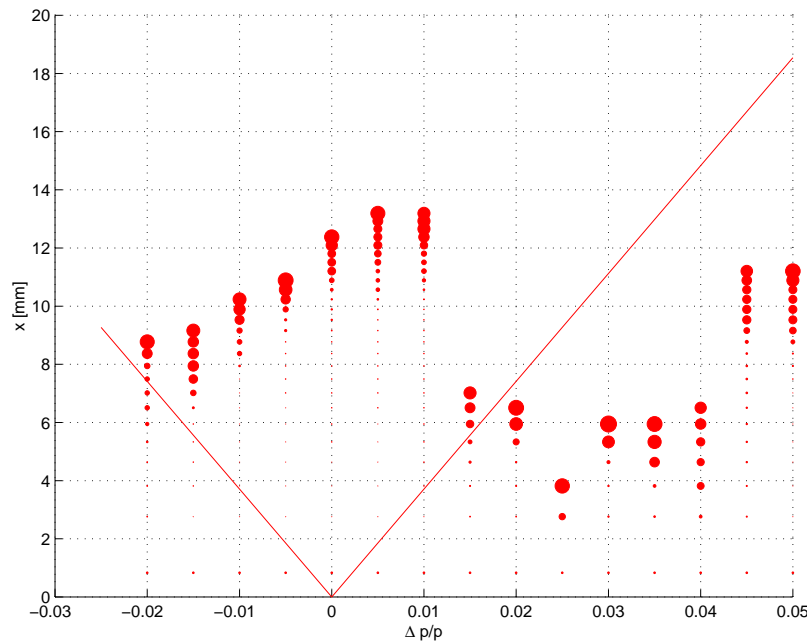
- structures in dynamic aperture can be identified with resonances
- chromaticity (even the linear one) is important

Off energy aperture studies/chromaticity



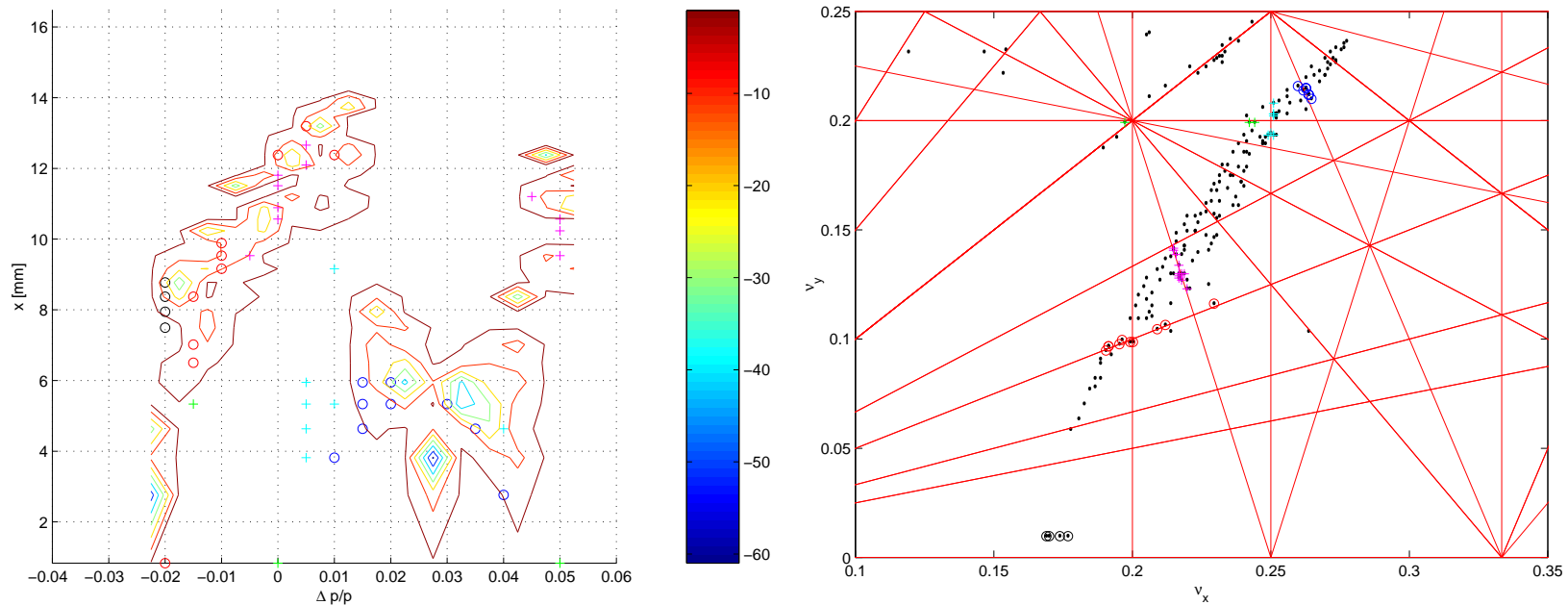
- with different linear chromaticity, off energy particles are influenced by different resonances \Rightarrow off energy dynamics can be very different
- effect can be fairly large

Off energy aperture studies/chromaticity II



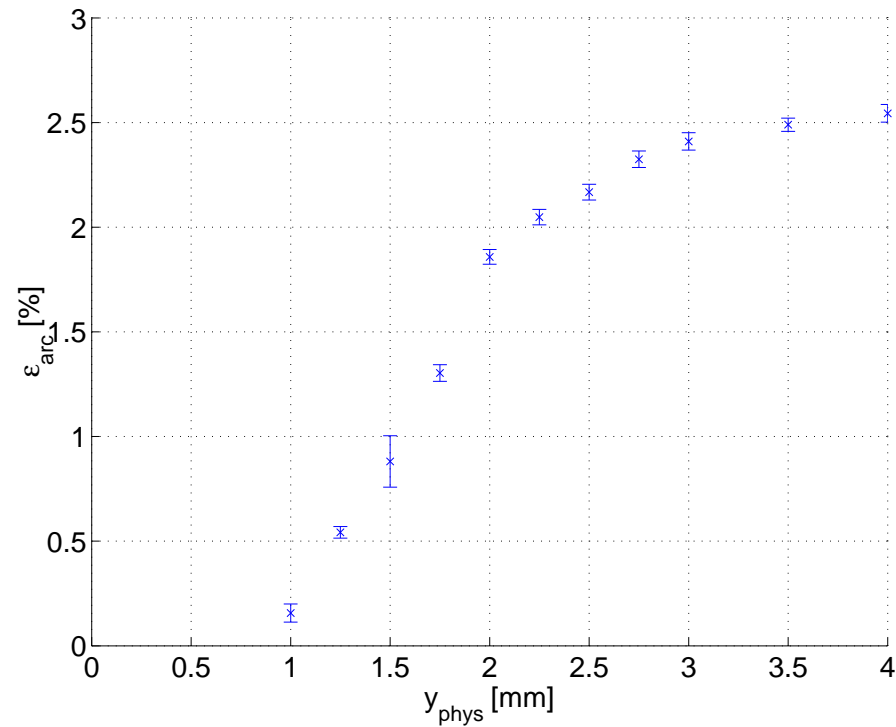
- due to linear chromaticity a region with strongly excited resonances is probed
- in current case they were especially strong because symmetry was broken (three-fold)

Off energy aperture studies/chromaticity III



- lossy **regions** in phase space can be easily identified with **resonances** (enhanced by periodicity breaking)
- explains why ALS lifetime is very sensitive to vertical chromaticity

Impact of Vertical Physical Aperture



- Vertical physical aperture has big impact on momentum aperture
- For ALS lifetime collapses at aperture of about 40-50 σ_y

Summary (Momentum Aperture)



- *Dynamic momentum aperture is important/dominant effect for Touschek lifetime*
 - ◆ Measurement method using **frequency analysis** provides a very powerful **model independent diagnostic tool**
 - ◆ Improvement possibilities limited in ALS (two sextupole families), but **large potential in newer light sources**

- *Agreement between the measurement and simulation is good*
 - ◆ \Rightarrow Confidence in **machine model** \Rightarrow can be used to **predict performance** of upgrades or new machines

- *Outlook:*
 - ◆ Reduction in measurement time
 - ◆ Apply transverse and longitudinal kick simultaneously