



Resonance driving term analysis

Studies on Lattice Calibration With Frequency Analysis of Betatron Motion

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Most of the material in this lecture from

Ricardo Bartolini (Diamond),

Frank Schmidt, R. Tomas (CERN) , ...

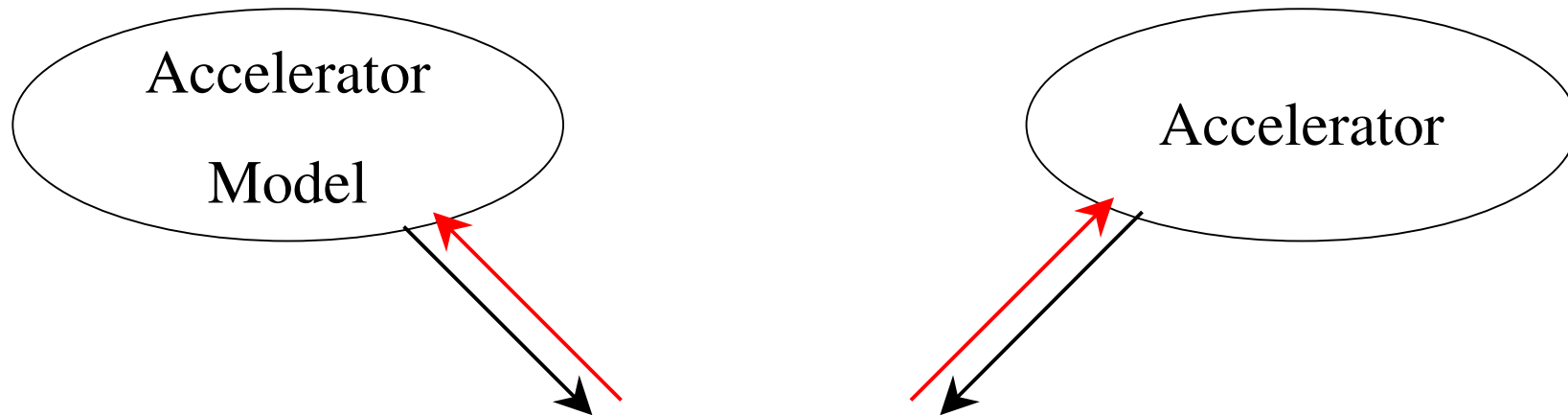
- Outline:

- ❖ Introduction

- NAFF
 - perturbative theory of betatron motion
 - SVD fit of lattice parameters

- ❖ DIAMOND Spectral Lines Analysis

- Linear Model
 - Nonlinear Model



- Closed Orbit Response Matrix (LOCO-like)
- Frequency Map Analysis
- Frequency Analysis of Betatron Motion (resonant driving terms)

Frequency Analysis of Betatron Motion and Lattice Model Reconstruction (1)



Accelerator Model



- tracking data at all BPMs
- spectral lines from model (NAFF)
- build a vector of Fourier coefficients

Accelerator



- beam data at all BPMs
- spectral lines from BPMs signals (NAFF)
- build a vector of Fourier coefficients

$$\bar{A} = (a_1^{(1)} \quad \dots \quad a_{NBPM}^{(1)} \quad \varphi_1^{(1)} \quad \dots \quad \varphi_{NBPM}^{(1)} \quad a_1^{(2)} \quad \dots \quad a_{NBPM}^{(2)} \quad \varphi_1^{(2)} \quad \dots \quad \varphi_{NBPM}^{(2)} \quad \dots)$$

Two arrows point from the text "build a vector of Fourier coefficients" in the previous blocks to the corresponding terms in the vector \bar{A} .

Define the distance between the two vector of Fourier coefficients

$$\chi^2 = \sum_k (A_{Model}(j) - A_{Measured}(j))^2$$

Frequency Analysis of Betatron Motion and Lattice Model Reconstruction (2)



Least Square Fit (SVD) of accelerator parameters θ
to minimize the distance χ^2 of the two Fourier coefficients vectors

- Compute the “Sensitivity Matrix” M
- Use SVD to invert the matrix M
- Get the fitted parameters

$$\Delta \bar{A} = M \bar{\theta}$$

$$M = U^T W V$$

$$\bar{\theta} = (V^T W^{-1} U) \Delta \bar{A}$$

MODEL \rightarrow TRACKING \rightarrow NAFF \rightarrow

Define the vector of Fourier Coefficients – Define the parameters to be fitted

SVD \rightarrow CALIBRATED MODEL

DIAMOND Layout



Main parameters:

100 MeV Linac

3 GeV Booster (158.4 m)

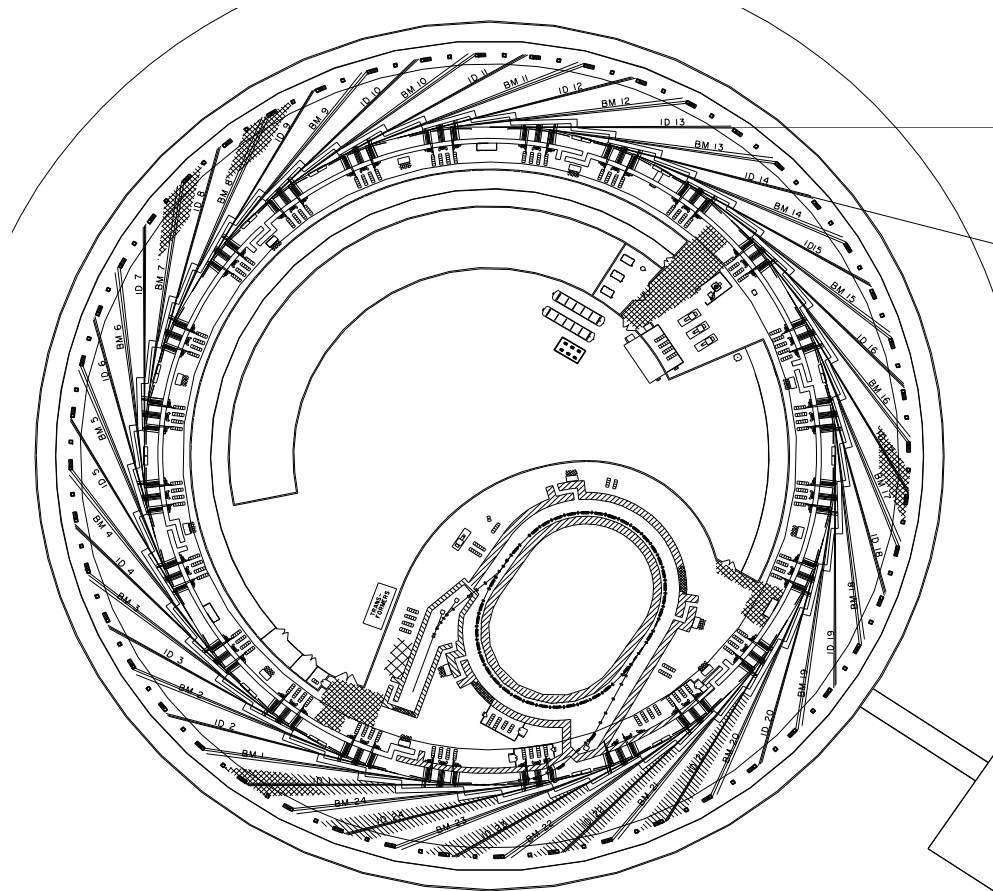
3 GeV Storage Ring (561.6 m)

24 cell DBA lattice

2 + 1 SC RF cavities

18 straight for ID (5 m)

6 long straights (8 m)



Commissioning end 2006

NAFF algorithm – J. Laskar (1988) (Numerical Analysis of Fundamental Frequencies)



Given the quasi-periodic time series of the particle orbit $(x(n); p_x(n))$,

- Find the main lines with the previous technique for tune measurement

⇒ ν_1 frequency, a_1 amplitude, ϕ_1 phase;

- build the harmonic time series

$$z_1(n) = a_1 e^{i\phi_1} e^{2\pi i \nu_1 n}$$

- subtract from the original signal
- analyze again the new signal $z(n) - z_1(n)$ obtained ●

The decomposition $z(n) = \sum_{k=1}^n a_k e^{i\phi_k} e^{2\pi i \nu_k n}$ **allows the**

Measurement of Resonant driving terms of non linear resonances

Frequency Analysis of Non Linear Betatron Motion

A.Ando (1984), J. Bengtsson (1988), R.Bartolini-F. Schmidt (1998)



The quasi periodic decomposition of the orbit

$$x(n) - ip_x(n) = \sum_{k=1}^n c_k e^{2\pi i \nu_k n} \quad c_k = a_k e^{i\phi_k}$$

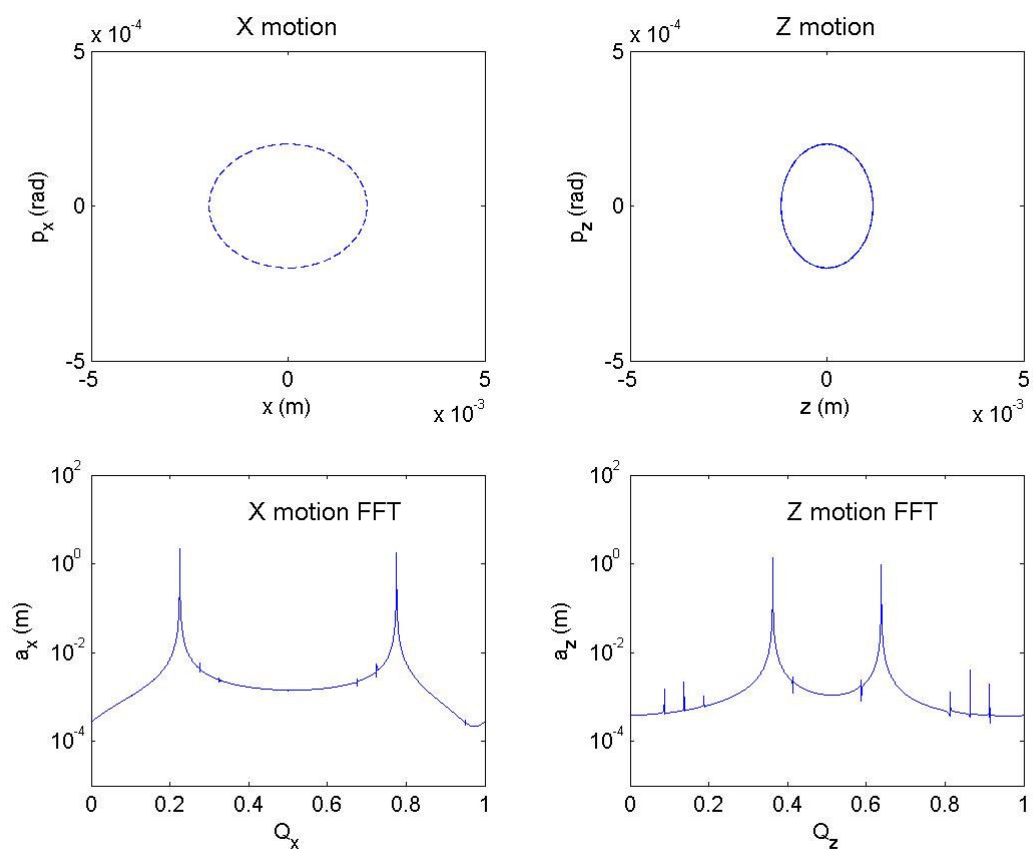
can be compared to the perturbative expansion of the non linear betatron motion

$$x(n) - ip_x(n) = \sqrt{2I_x} e^{i(2\pi Q_x n + \psi_0)} + \\ - 2i \sum_{jklm} j s_{jklm} (2I_x)^{\frac{j+k-1}{2}} (2I_y)^{\frac{l+m}{2}} e^{i[(1-j+k)(2\pi Q_x n + \psi_{x0}) + (m-l)(2\pi Q_y n + \psi_{y0})]}$$

Each resonance driving term s_{jklm} contributes to the Fourier coefficient of a well precise spectral line

$$\nu(s_{jklm}) = (1 - j + k)Q_x + (m - l)Q_y$$

Spectral Lines for DIAMOND low emittance lattice (.2 mrad kick in both planes)

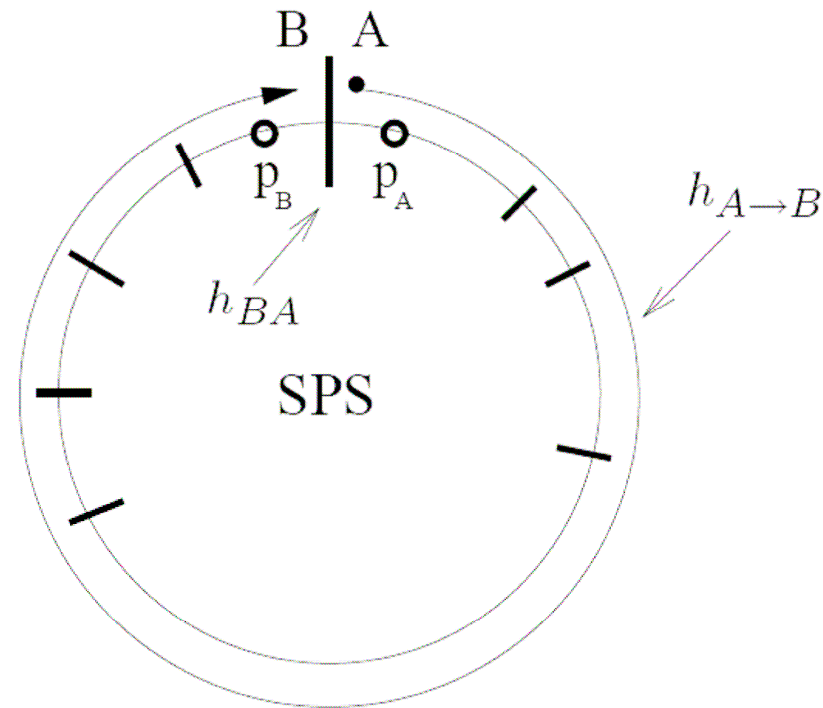


Spectral Lines detected with NAFF algorithm

e.g. Horizontal:

- (1, 0) $1.10 \cdot 10^{-3}$ horizontal tune
- (0, 2) $1.04 \cdot 10^{-6}$ $Q_x - 2 Q_z$
- (-3, 0) $2.21 \cdot 10^{-7}$ $4 Q_x$
- (-1, 2) $1.31 \cdot 10^{-7}$ $2 Q_x + 2 Q_z$
- (-2, 0) $9.90 \cdot 10^{-8}$ $3 Q_x$
- (-1, 4) $2.08 \cdot 10^{-8}$ $2 Q_x + 4 Q_z$

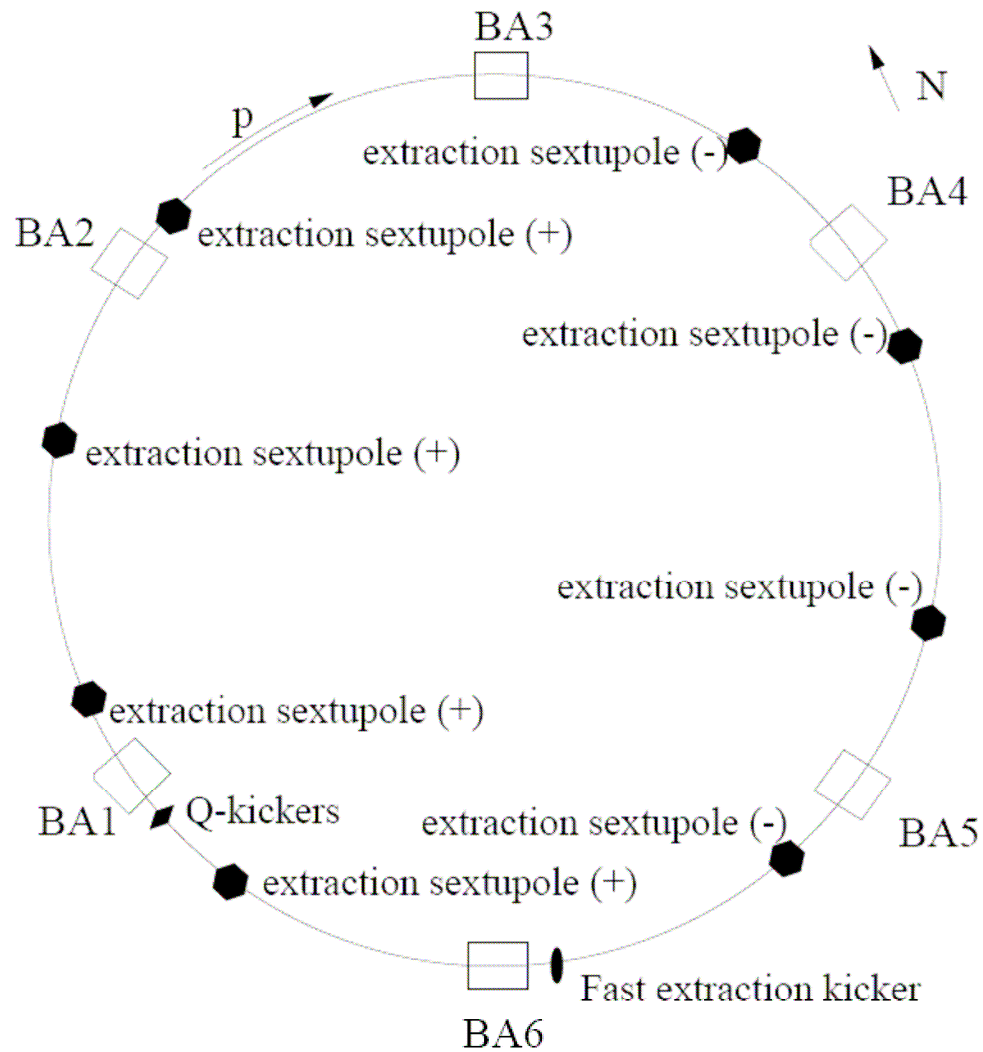
Longitudinal Variation of Driving Terms



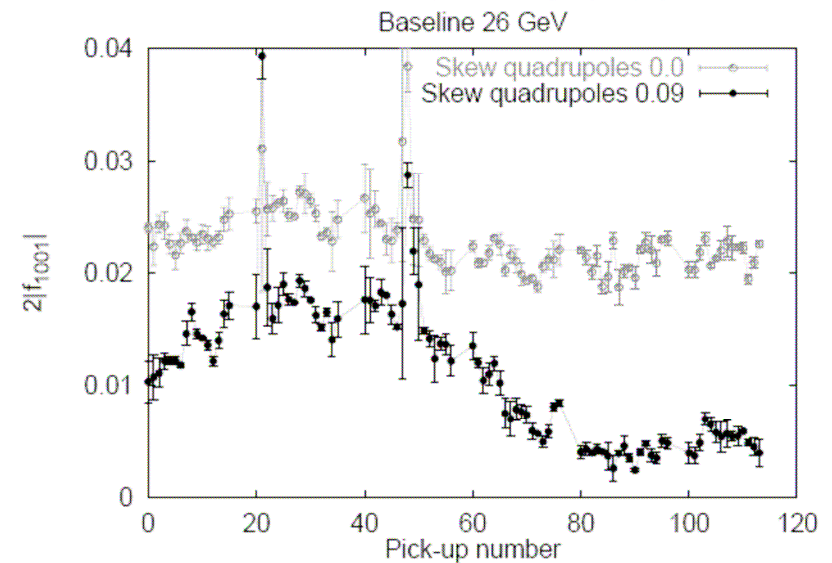
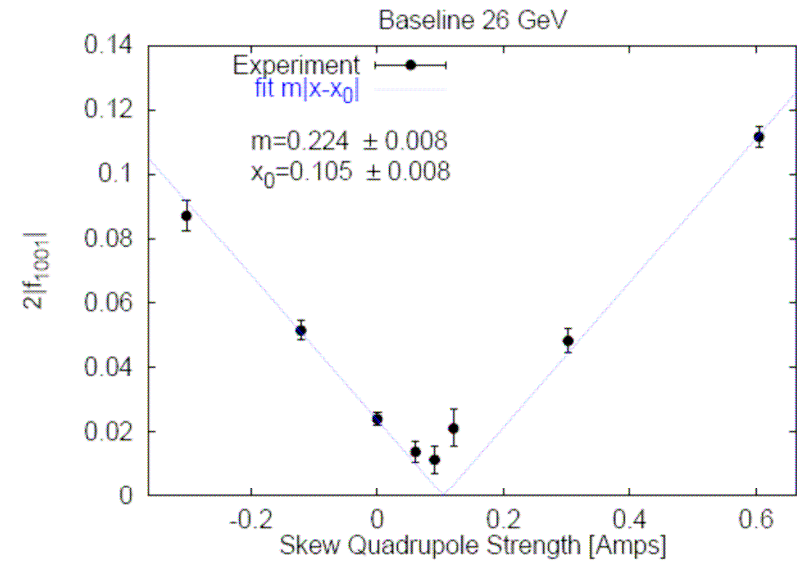
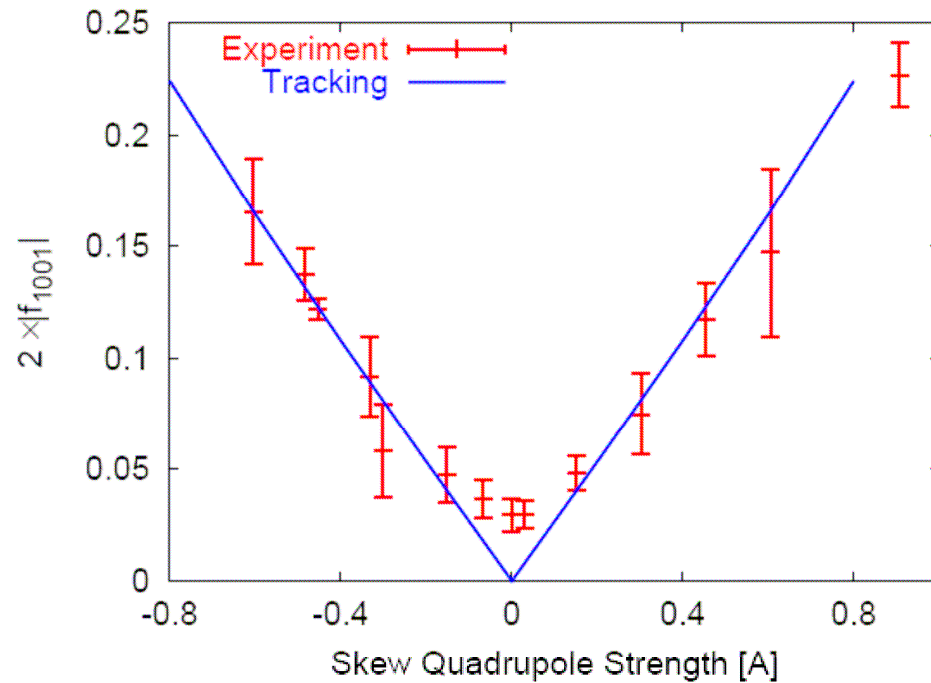
$$\left. \begin{aligned} h_A &= h_{A \rightarrow B} \circ h_{B \rightarrow A} \\ h_B &= h_{B \rightarrow A} \circ h_{A \rightarrow B} \end{aligned} \right\} h_A \neq h_B$$

- \Rightarrow Amplitude of Driving Terms not the same
- \Rightarrow Amplitude constant between lattice errors

Early (2000-2002) Work: SPS experiments

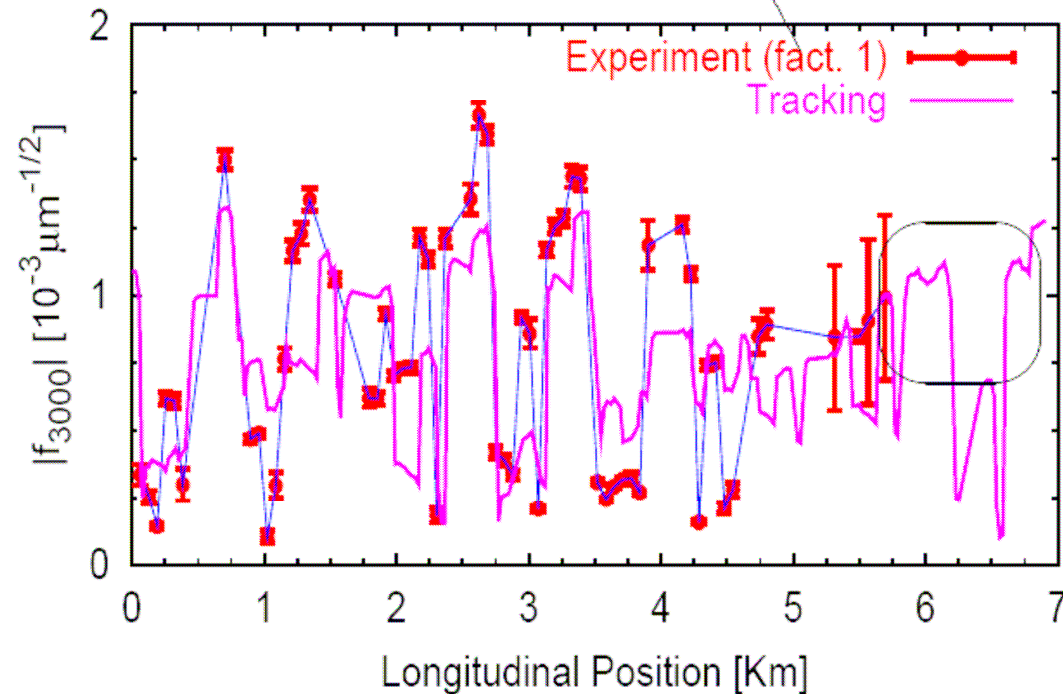


Linear Coupling in SPS



Sextupole Driving Terms

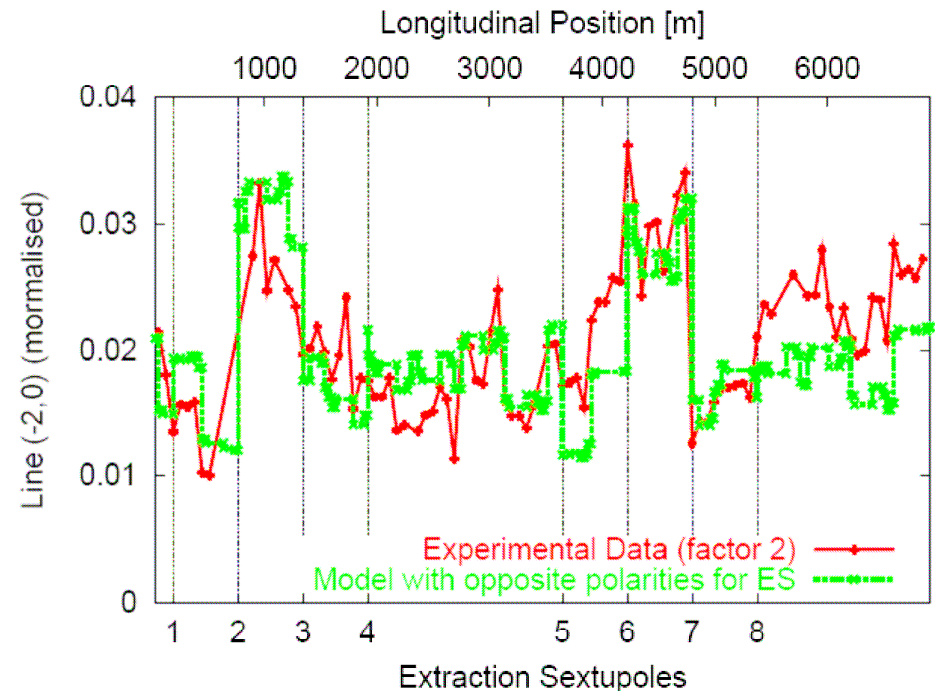
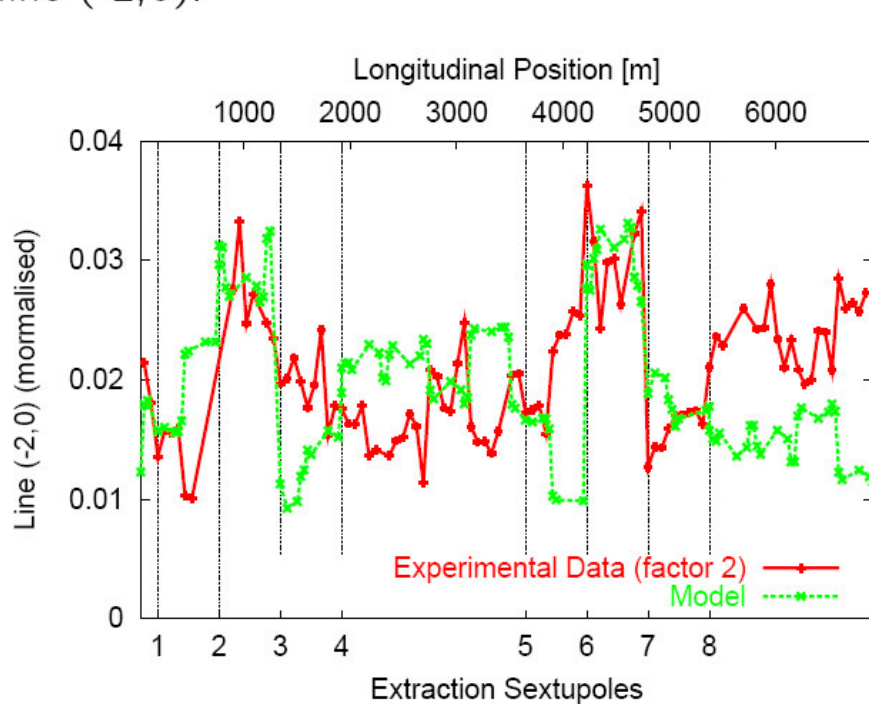
Bare SPS, $Q_x=26.62$ (No decoherence).



⇒ Local discrepancies may be due to unknown sextupole sources.

Sextupole Driving Terms with Extraction Sext.

The resonance (3,0) introduces the spectral Change polarities of the extraction sextupoles? line (-2,0).



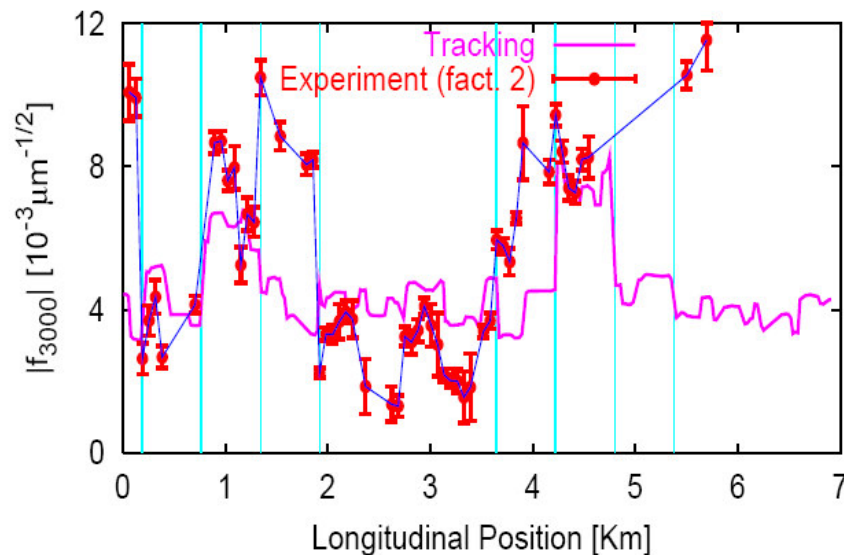
⇒ We have a problem!

Hardware checks confirmed that these sextupoles had opposite polarities.

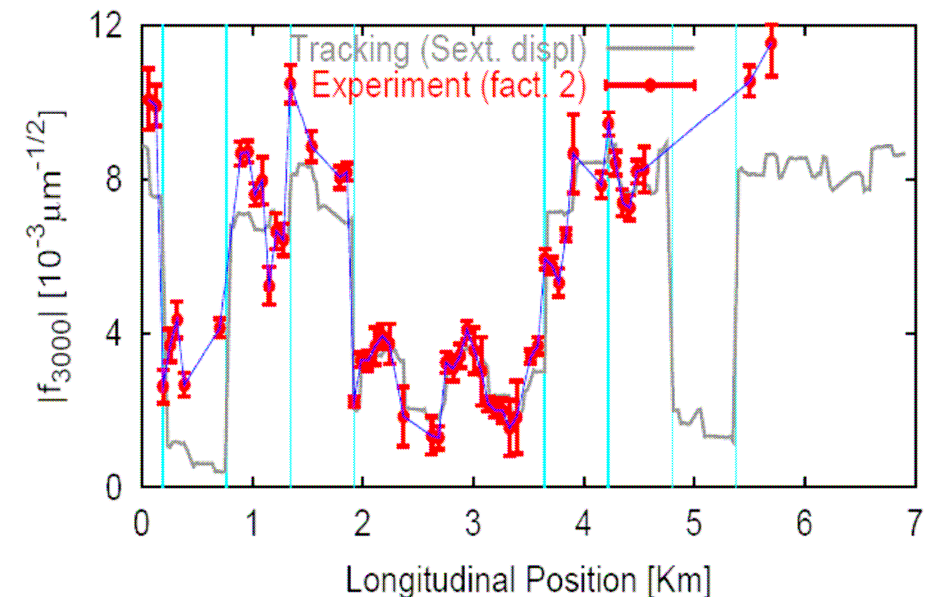
Sextupole Driving Terms with Extraction Sext.

Extraction Sextupoles powered to
 + + + + - - - - 30 A.
 Data fully decohered \Rightarrow Line reduced by a factor 2.

The closed orbit as measured from pick-ups is introduced at the extraction sextupoles in the model.



\Rightarrow Large discrepancies: We have a problem!

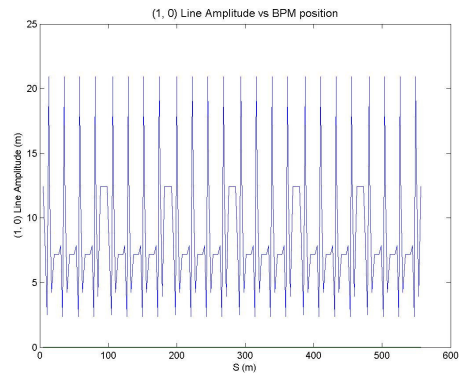
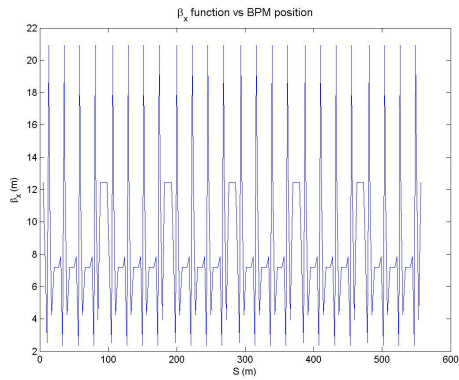


\Rightarrow Improvement due to beta-beating and phase differences in the model.

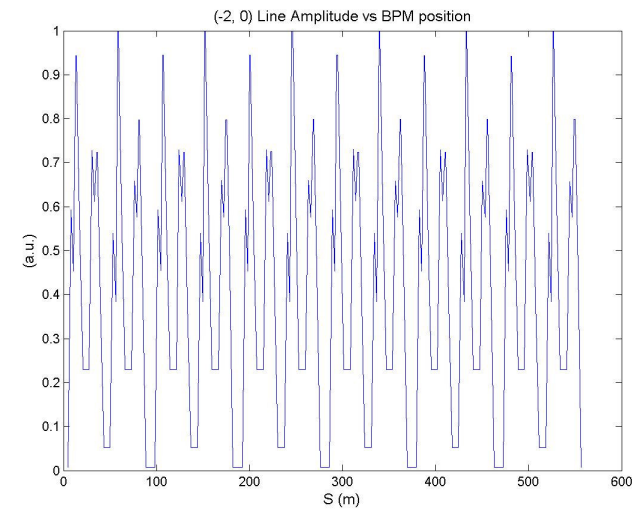
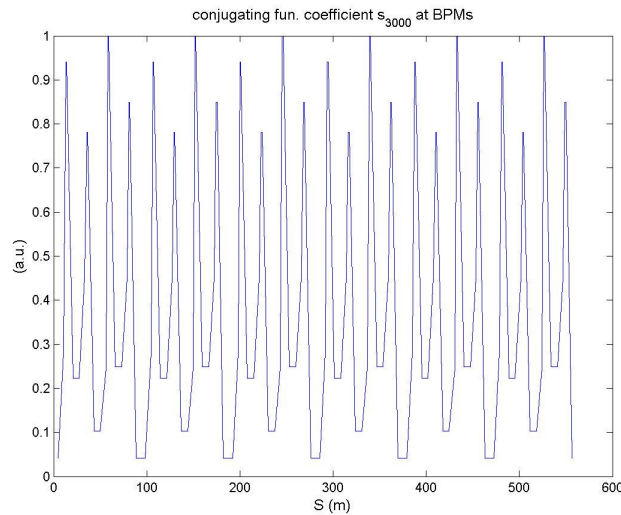
Amplitude of Spectral Lines for low emittance DIAMOND lattice computed at all the BPMs



Main spectral line (Tune Q_x)



(-2, 0) spectral line: resonance driving term h_{3000} ($3Q_x = p$) at all BPMs



- The amplitude of the tune spectral line replicates the β functions
- The amplitude of the $(-2, 0)$ show that third order resonance is well compensated within one superperiod. Some residual is left every two cells ($5\pi/2$ phase advance)

Example: DIAMOND with random misalignments (100 μm r.m.s) in chromatic sextupoles to generate linear coupling



The coupled linear motion in each plane can be written in terms of the coupling matrix

$$C = \begin{pmatrix} c_{11} & c_{12} \\ c_{21} & c_{22} \end{pmatrix}$$

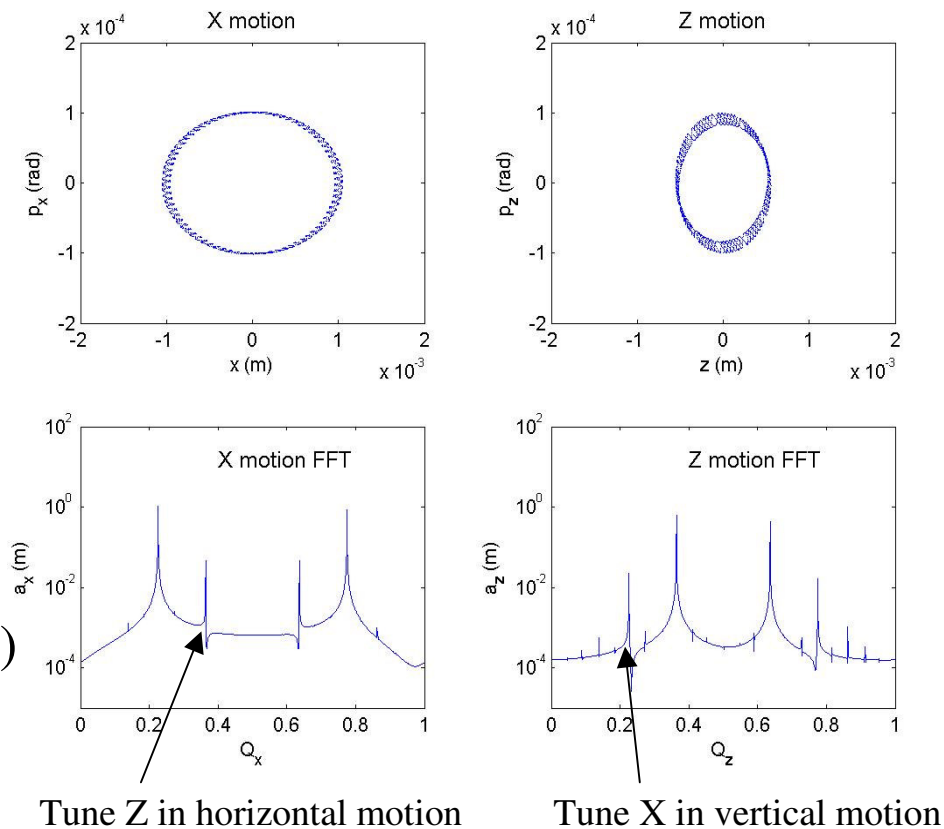
e.g. for the horizontal motion

$$\zeta_x = x - ip_x = a_1 e^{i(\phi_u + \delta_u)} + a_2 e^{-i(\phi_u + \delta_u)} + a_3 e^{i(\phi_v + \delta_v)} + a_4 e^{-i(\phi_v + \delta_v)}$$

a_3 and a_4 depend linearly on c_{ij}

- two frequencies (the H tune and V tune)
- no detuning with amplitude

0.1 mrad kick in both planes



(0,1) spectral line for low emittance DIAMOND lattice computed at all the BPMs (V misalignment errors added to chromatic sextupoles)

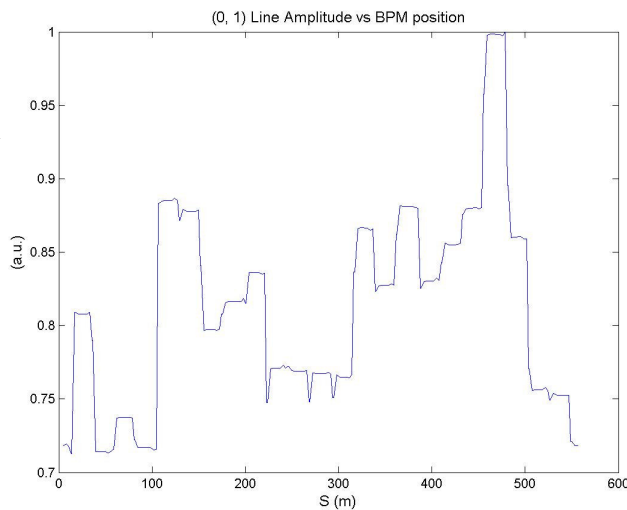


$$h_{1001p} = \frac{1}{2\pi} \int_0^{2\pi R} a_2(s) \left(\frac{\beta_x}{2}\right)^{1/2} \left(\frac{\beta_z}{2}\right)^{1/2} e^{i(W_x(s)-W_z(s))-i\frac{ps}{R}} ds$$

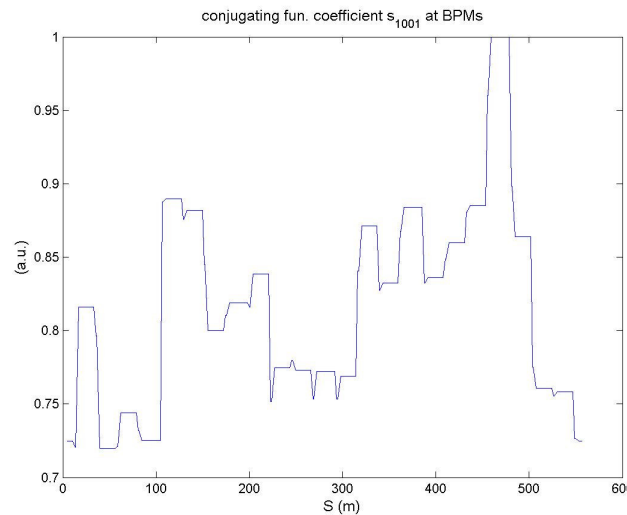
$$s_{1001p} = \frac{ih_{1001p}}{Q_x - Q_z + p} \quad s_{1001}(s) = \frac{1}{2\pi R} \sum_{p=-\infty}^{\infty} s_{1001p} e^{ips/R}$$

The resonance driving term h_{1001} contributes to the (0, 1) spectral line in horizontal motion

from analytical formula



from spectral line



The amplitude of the (0, 1) spectral line replicates well the s dependence of the difference resonance $Q_x - Q_z$ driving term

DIAMOND Spectral Lines Analysis

- ❖ Horizontal Misalignment of sextupoles (β – beating)
- ❖ Vertical Misalignment of sextupoles (linear coupling)
- ❖ Gradient errors in sextupoles (non linear resonances)

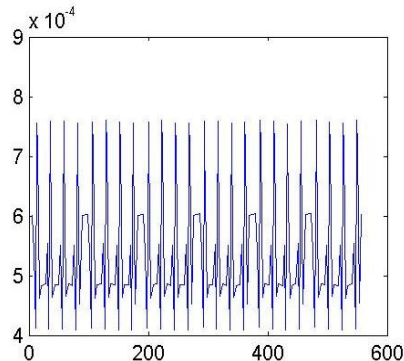
Horizontal misalignment of a set of 24 sextupoles with 100 μm rms (β - beating correction)



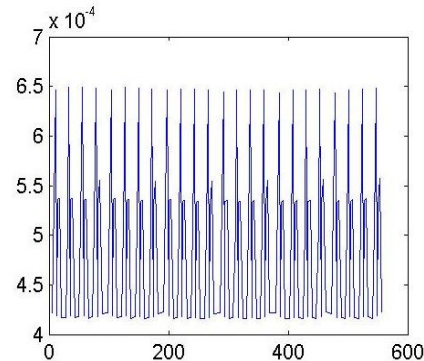
The generated normal quadrupole components introduce a β - beating.

- we build the vector of Fourier coefficients of the horizontal and vertical tune line
- we use the horizontal misalignments as fit parameters

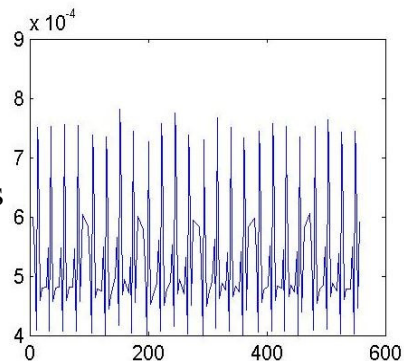
H tune line
(no misalignments)



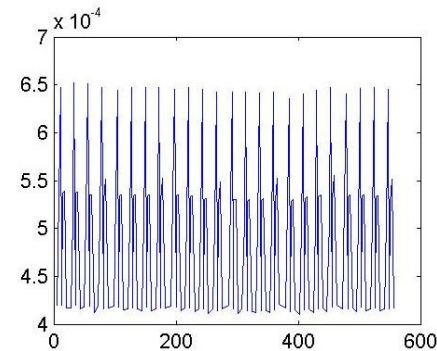
V tune line
(no misalignments)



H tune line
with misalignments



V tune line
with misalignments



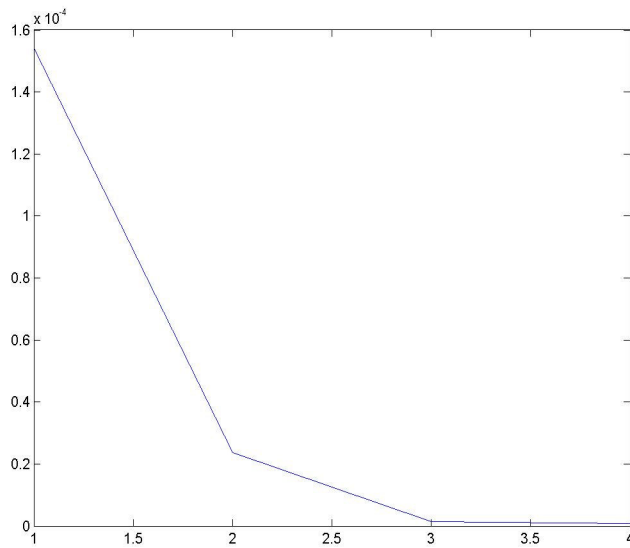
SVD on sextupoles horizontal misalignments



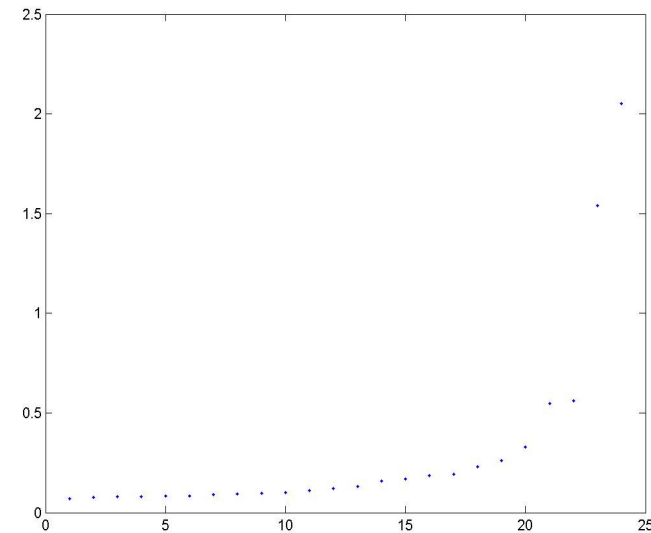
We build the vector $\bar{A} = (a_1^{H(1,0)} \quad \dots \quad a_{NBPM}^{H(1,0)} \quad a_1^{V(0,1)} \quad \dots \quad a_{NBPM}^{V(0,1)})$

containing the amplitude of the tune lines in the two planes at all BPMs

We minimize the sum $\chi^2 = \sum_j (A_{Model}(j) - A_{Measured}(j))^2$



χ^2 as a function of the iteration number



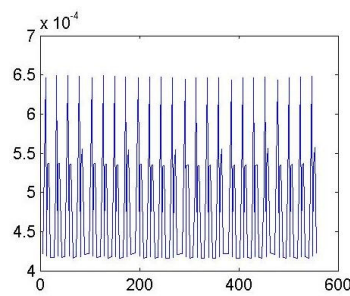
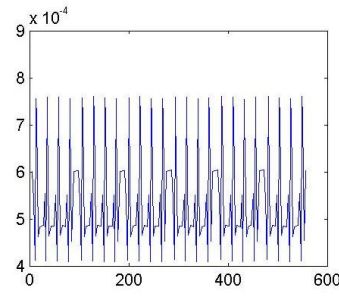
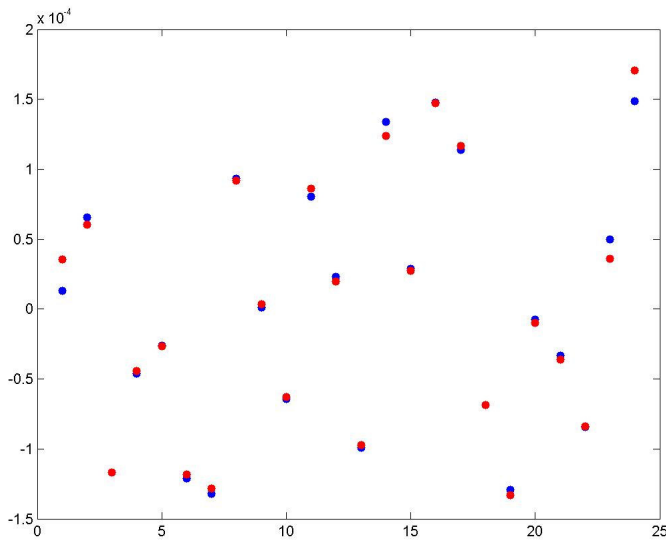
Example of SVD principal values

Fitted values for the 24 horizontal sextupole misalignments obtained from the SVD

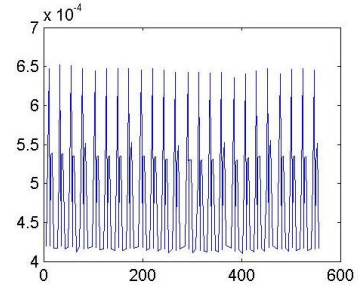
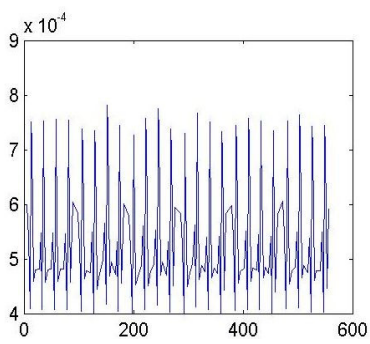


Blue dots = assigned misalignments

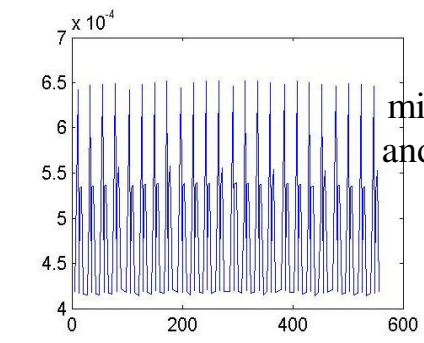
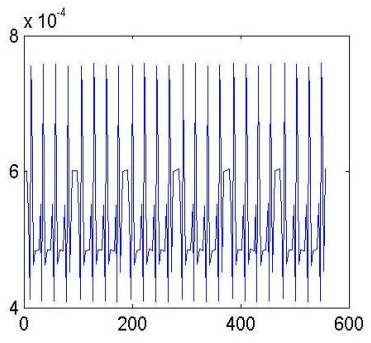
Red dots = reconstructed misalignments



no misalignments



with misalignments

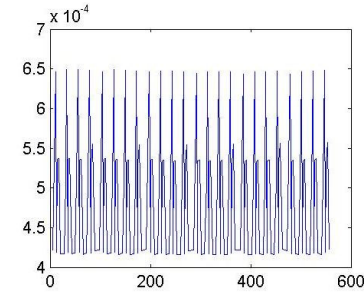
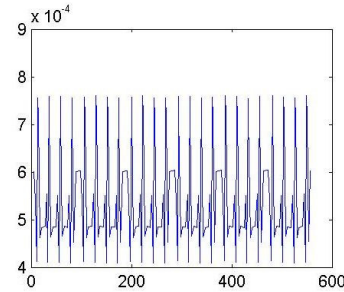
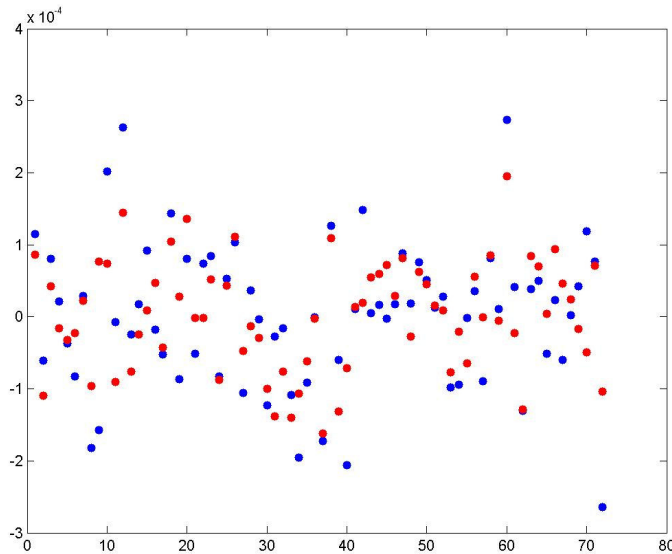


with misalignments and corrections

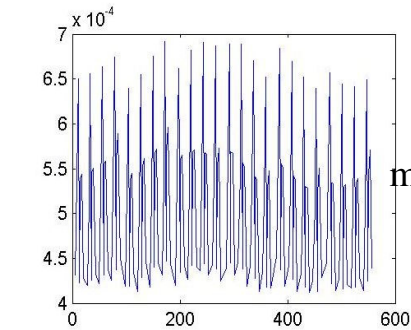
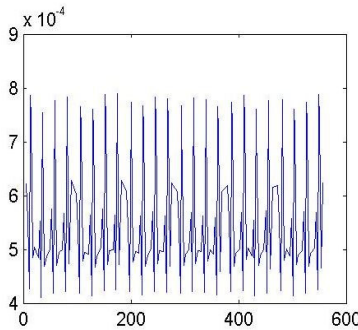
Fitted values for the 72 horizontal sextupole misalignments obtained from the SVD



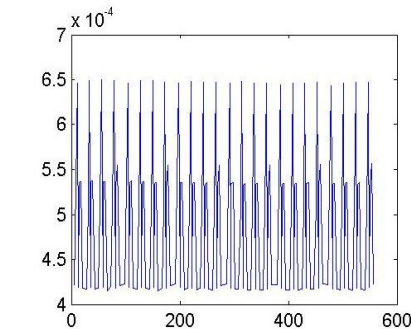
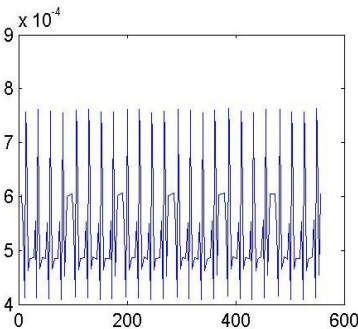
Blu dots = assigned misalignments
Red dots = reconstructed misalignments



no
misalignments



with
misalignments



with
alignments
& corrections

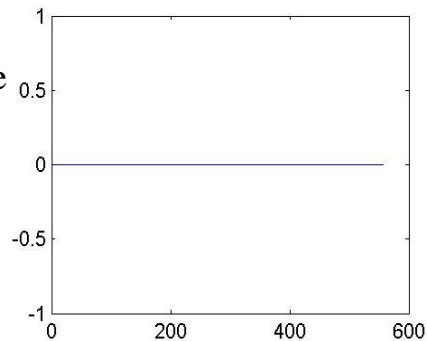
Vertical misalignment of a set of 24 sextupoles with 100 μm rms (linear coupling correction)



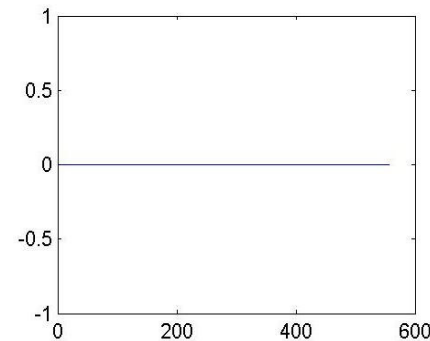
The generated skew quadrupole components introduce a linear coupling.

- we build the vector of Fourier coefficients of the (0, 1) line in the H plane
- we use the vertical misalignments as fit parameters

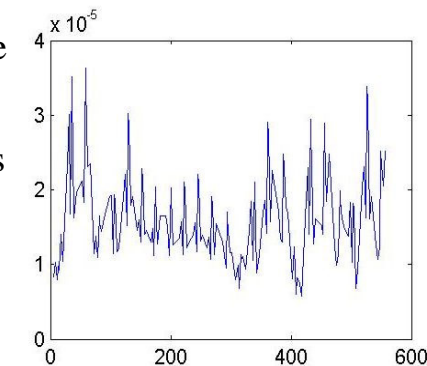
(0,1) line amplitude
in H plane
(no misalignments)



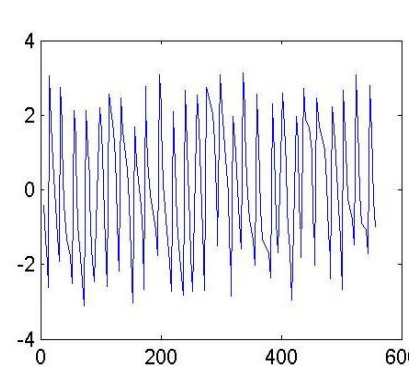
(0,1) line phase in H
plane
(no misalignments)



(0,1) line amplitude
in H plane
with misalignments



(0,1) line phase in H
plane
with misalignments



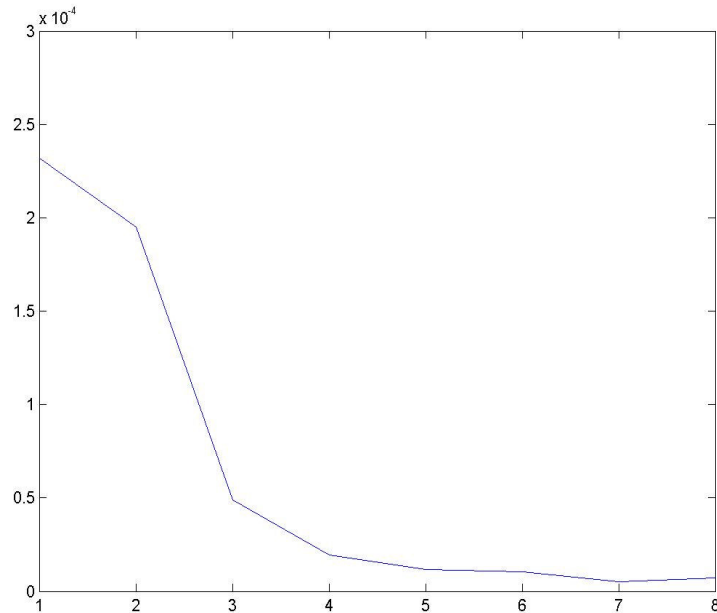
SVD on sextupole vertical misalignments



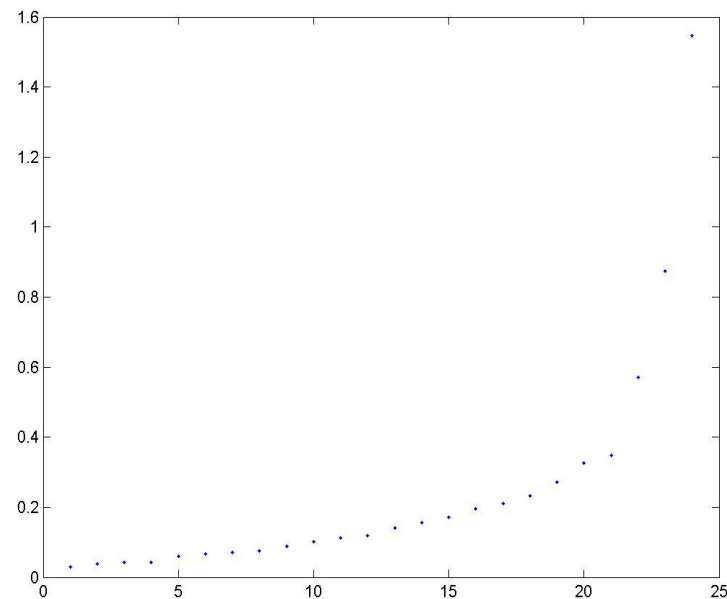
We build the vector $\bar{A} = (a_1^{H(0,1)} \quad \dots \quad a_{NBPM}^{H(0,1)} \quad \phi_1^{H(0,1)} \quad \dots \quad \phi_{NBPM}^{H(0,1)})$

containing the amplitude and phase of the (0, 1) line in the H planes at all BPMs

We minimize the sum $\chi^2 = \sum_j (A_{Model}(j) - A_{Measured}(j))^2$



χ^2 as a function of the iteration number



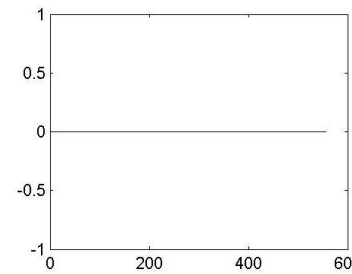
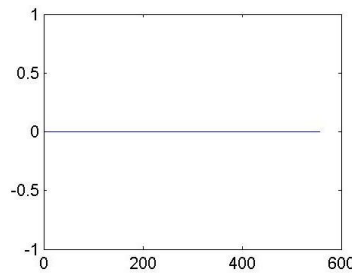
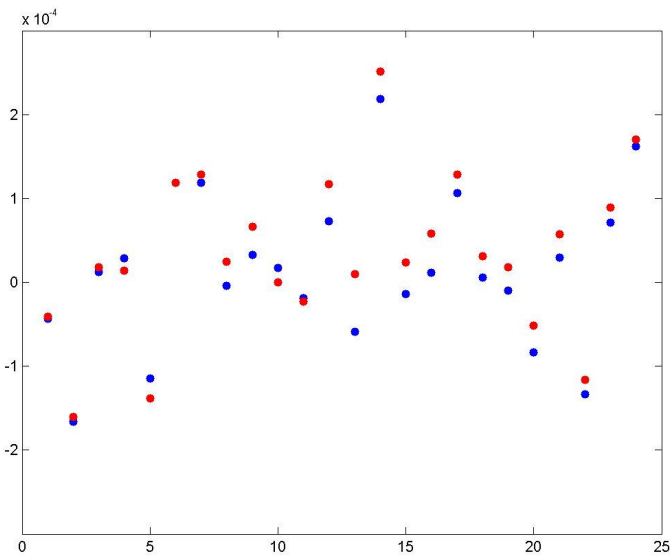
Example of SVD principal values

Fitted values for the 24 vertical sextupole misalignments obtained from SVD

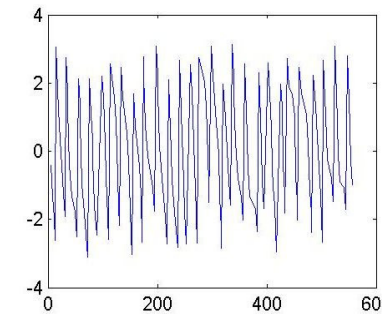
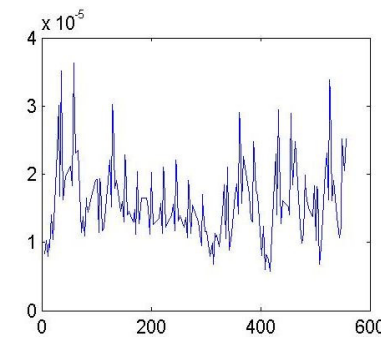


Blu dots = assigned misalignments

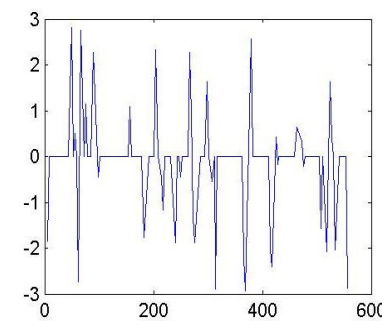
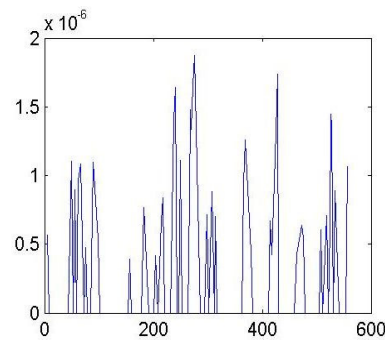
Red dots = reconstructed misalignments



no misalignments



with misalignments



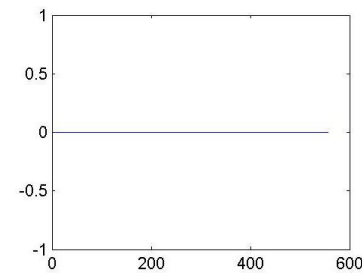
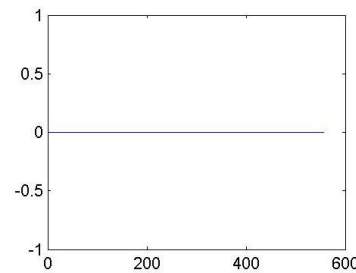
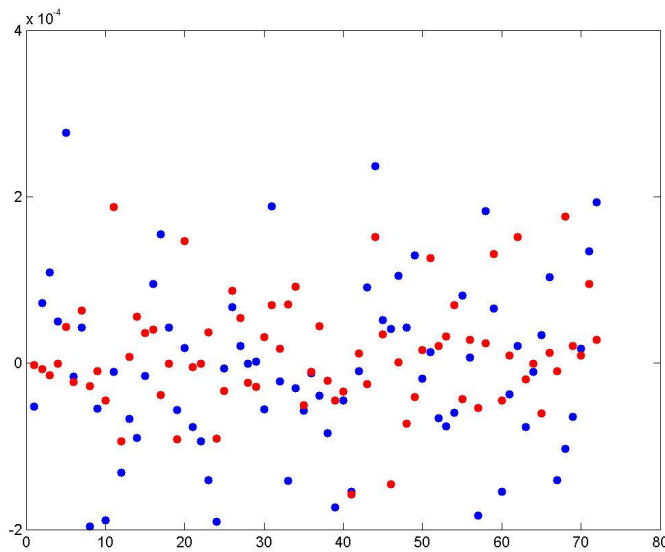
with alignments corrections

Fitted values for the 72 vertical sextupole misalignments obtained from SVD

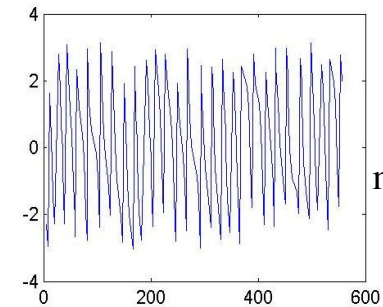
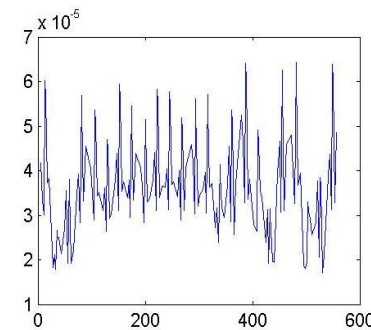


Blue dots = assigned misalignments

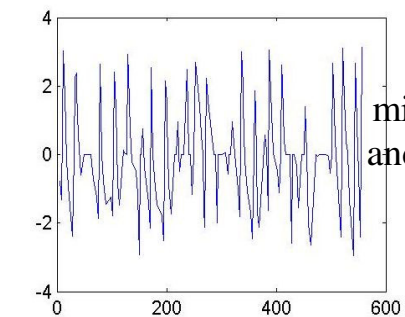
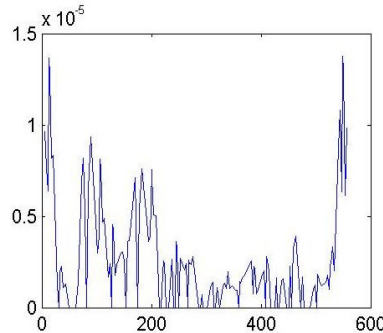
Red dots = reconstructed misalignments



no misalignments



with misalignments



with misalignments and corrections

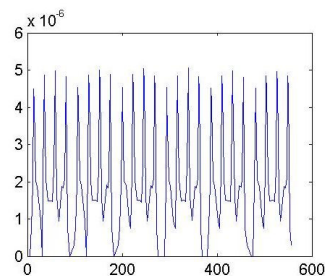
Sextupoles gradient errors applied to 24 sextupoles ($dK_2/K_2 = 5\%$)



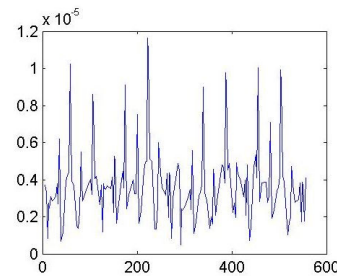
The sextupole gradient errors spoil the compensation of the third order resonances, e.g $3Q_x = p$ and $Q_x - 2Q_z = p$

- we build the vector of Fourier coefficients of the H(-2,0) and H(0,2) line
- we use the errors gradients as fit parameters

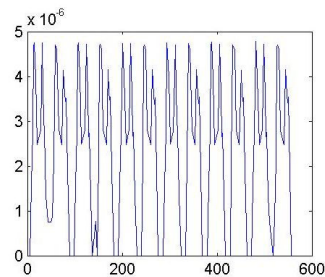
(0,2) line amplitude
in H plane
(no gradient errors)



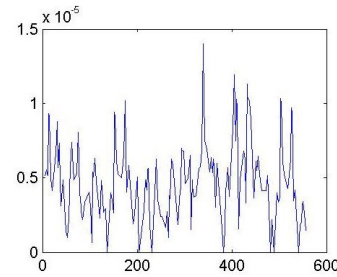
(0,2) line amplitude
in H plane
with gradient errors



(-2,0) line amplitude
in H plane
(no gradient errors)



(-2,0) line amplitude
in H plane
with gradient errors



SVD on sextupole gradient errors



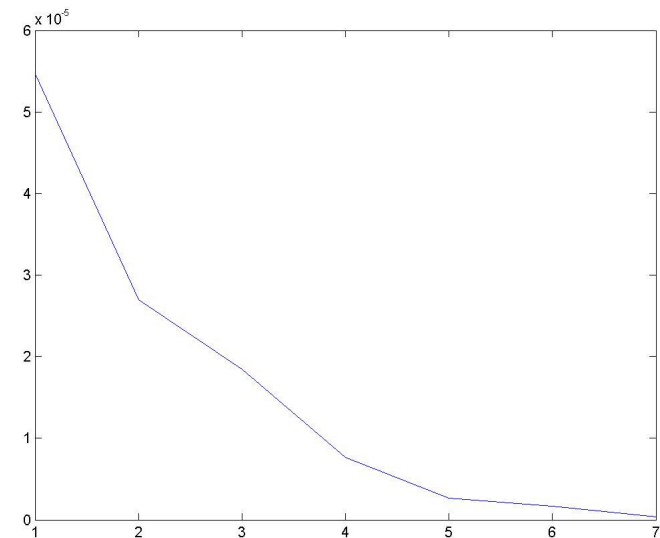
We build the vector $\bar{A} = \left(a_1^{H(-2,0)} \dots a_{NBPM}^{H(-2,0)} a_1^{H(0,2)} \dots a_{NBPM}^{H(0,2)} \right)$

containing the amplitudes at all BPMs

- the $(-2, 0)$ line in the H plane related to h_{3000}
- the $(0, 2)$ line in the H plane related to h_{1002}

We minimize the sum

$$\chi^2 = \sum_j \left(A_{Model}(j) - A_{Measured}(j) \right)^2$$



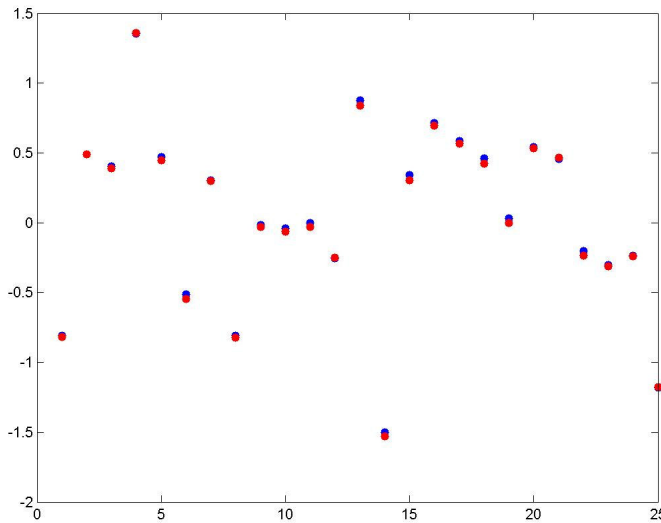
χ^2 as a function of the iteration number

Fitted values for the 24 sextupoles gradients errors obtained from SVD

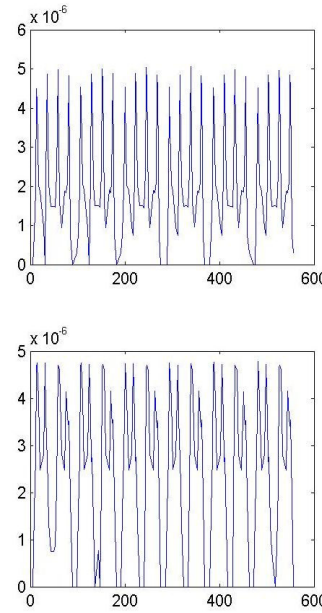


Blu dots = assigned misalignments

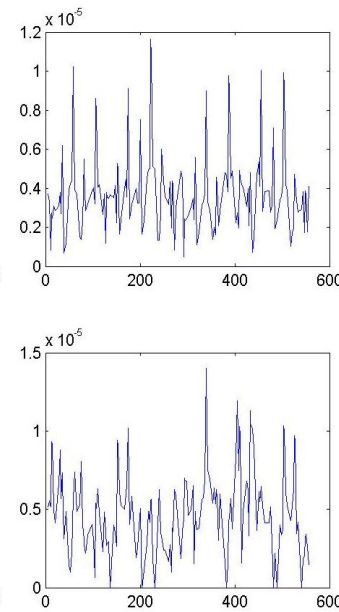
Red dots = reconstructed misalignments



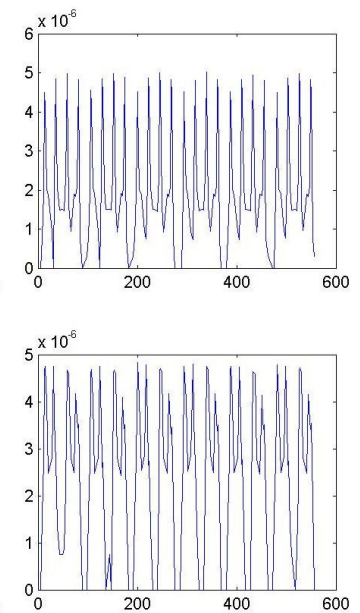
no
gradient errors



with
gradient errors



with gradient
errors and
corrections

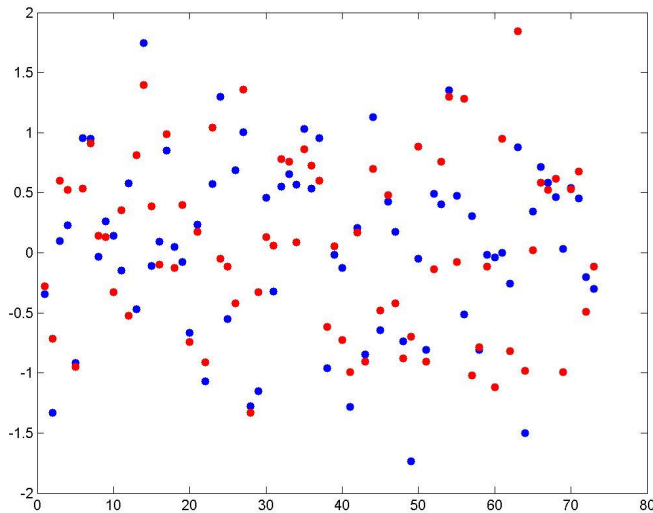


Fitted values for the 72 sextupoles gradients errors obtained from SVD

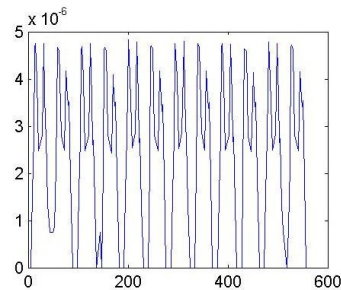
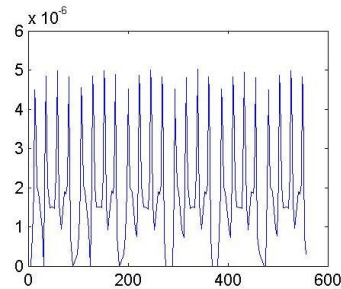


Blu dots = assigned misalignments

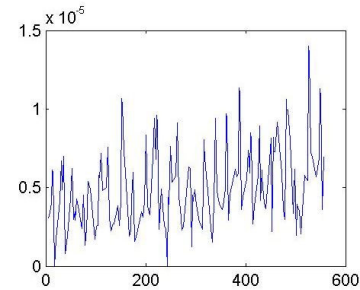
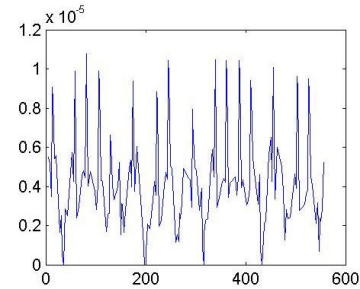
Red dots = reconstructed misalignments



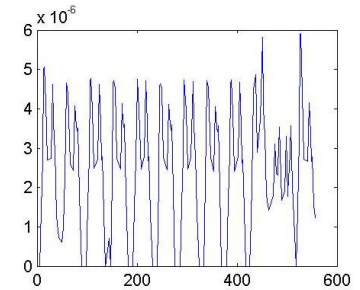
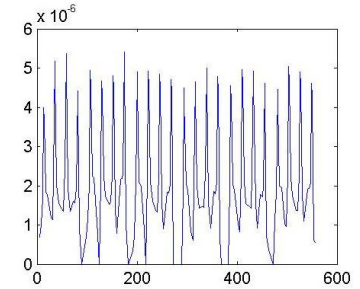
no
gradient errors



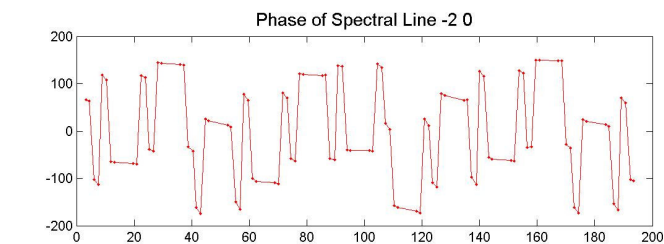
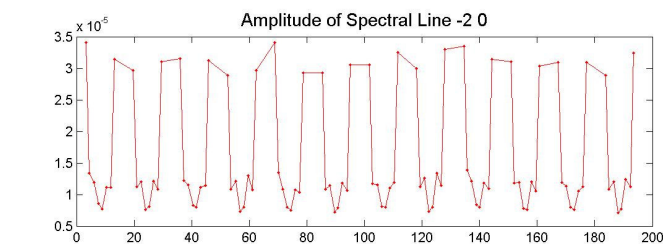
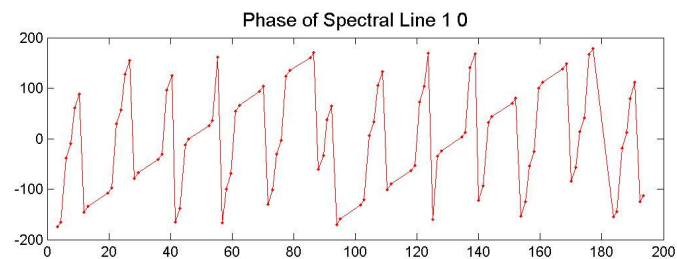
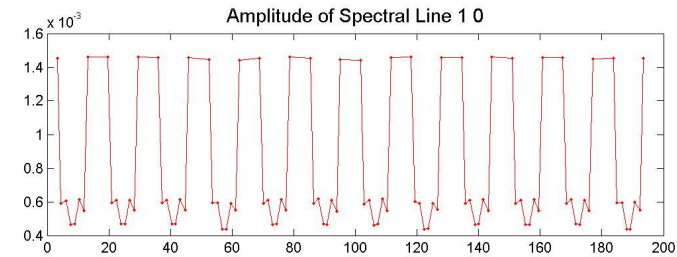
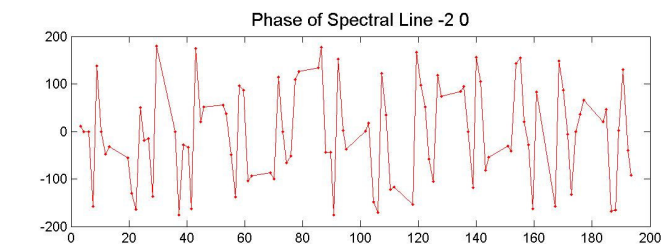
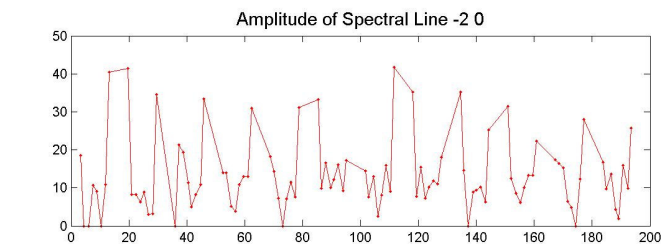
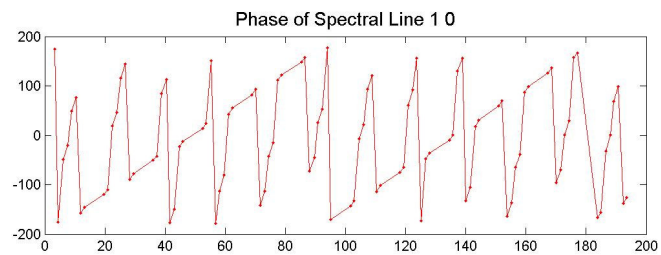
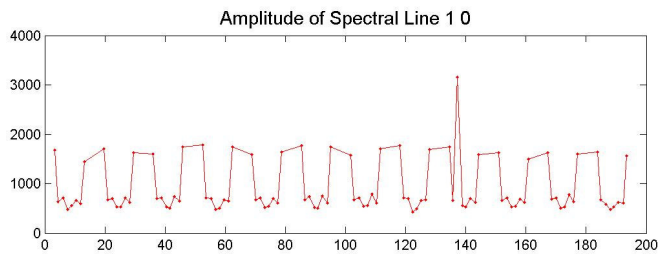
with
gradient errors



with gradient
errors and
corrections



ALS example (very early results)



Summary



- **Resonance driving term analysis provides quantitative information about nonlinearities in the machine**
- **It allows to measure the local distribution of the dominant nonlinearities**
- **However, it does not give a direct information about how harmful the nonlinearities are**
- **Theoretically it can provide a method similar to orbit response matrix analysis (or phase advance, ...) to measure not just the gradient and skew gradient distribution, but also the setxupole, (octupole), ... How well this will work experimentally is not quite clear, yet:**
 - **Can we use the spectral lines to recover the LINEAR and NON LINEAR machine model with a Least Square method?**
- **SPS with a few very large nonlinearities worked well. Diamond simulations look encouraging. ALS measurements so far seem BPM resolution limited.**



Further Reading

- ❖ R. Bartolini, et al. 'Measurement of Resonance Driving Terms by Turn-by-Turn Data', Proceedings of PAC 1999, 1557, New York (1999)
- ❖ W. Fischer, et al. 'Measurement of Sextupolar Resonance Driving Terms in RHIC', Proceedings of PAC 2003, Portland
- ❖ R. Bartolini and F. Schmidt, LHC Project note 132, Part. Accelerators. 59, pp. 93-106, (1998).
- ❖ F. Schmidt, R. Tom´as, A. Faus–Golfe, CERN-SL-2001-039-AP Geneva, CERN and IPAC 2001.
- ❖ M. Hayes, F. Schmidt, R. Tom´as, EPAC 2002 and CERN-SL-2002-039-AP Geneva, 25 Jun 2002.
- ❖ F. Schmidt, CERN SL/94-56 (AP) Update March 2000.
- ❖ R. Bartolini and F. Schmidt, CERN SL–Note–98–017 (AP).